BOUCHAUD'S MODEL EXHIBITS TWO DIFFERENT AGING REGIMES IN DIMENSION ONE

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ABSTRACT. Let E_i be a collection of i.i.d. exponential random variables. Bouchaud's model on $\mathbb Z$ is a Markov chain X(t) whose transition rates are given by $w_{ij} = \nu \exp(-\beta((1-a)E_i - aE_j))$ if i, j are neighbours in $\mathbb Z$. We study the behaviour of two correlation functions: $\mathbb P[X(t_w+t)=X(t_w)]$ and $\mathbb P[X(t')=X(t_w)\forall t'\in[t_w,t_w+t]]$. We prove the (sub)aging behaviour of these functions when $\beta>1$ and $a\in[0,1]$.

1. Introduction

Aging is an out-of-equilibrium physical phenomenon that is gaining considerable interest in contemporary physics and mathematics. An extensive literature exists in physics (see Bouchaud, Cugliandolo, Kurchan, and Mézard (1998) and their references). The mathematical literature is substantially smaller, although some progress was achieved in recent years (Ben Arous, Dembo, and Guionnet (2001); Ben Arous, Bovier, and Gayrard (2002a); Dembo, Guionnet, and Zeitouni (2001); Fontes, Isopi, and Newman (2002), see also Ben Arous (2002) for a survey).

The following model has been proposed by Bouchaud as a toy model for studying the aging phenomenon. Let $G = (\mathcal{V}, \mathcal{E})$ be a graph, $E = \{E_i\}_{i \in \mathcal{V}}$ be the collection of i.i.d. random variables indexed by vertices of this graph with the common exponential distribution with mean one. We consider the continuous time Markov chain X(t) with state space \mathcal{V} , such that

$$\mathbb{P}(X(t+dt)=j|X(t)=i,E) = \begin{cases} w_{ij}dt & \text{if } i, j \text{ are connected in } G\\ 0 & \text{otherwise.} \end{cases}$$
(1)

The transition rates w_{ij} are defined by

$$w_{ij} = \nu \exp\left(-\beta((1-a)E_i - aE_j)\right). \tag{2}$$

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The parameter β denotes, as usually, the inverse temperature and the parameter $a, 0 \le a \le 1$, drives the "symmetry" of the model. The value of ν fixes the time scale and is irrelevant for our paper, we thus set $\nu = 1$.

This model has been studied when G is \mathbb{Z} and a=0 by Fontes, Isopi, and Newman (1999, 2002). It is an elementary model when G is the complete graph, which is a good ansatz for the dynamics of the REM (see Ben Arous, Bovier, and Gayrard (2002b)).

The time spent by the system at site i grows with the value of E_i . The value of E_i can thus be regarded as the depth of the trap at the site i. The model is sometimes referred to as "Bouchaud's trap model." It describes the motion of the physical system between the states with energies $-E_i$. It can be regarded as a useful rough approximation of spin-glass dynamics. The states of Bouchaud's trap model correspond to a subset of all possible states of the spin-glass system with exceptionally low energy. This justifies in a certain sense the exponential distribution of E_i since it is the distribution of extreme values. The idea behind this model is that the spin-glass dynamics spends most of the time in the deepest states and it passes through all others extremely quickly. Thus, only the extremal states are important for the long time behaviour of dynamics, which justifies formally the introduction of Bouchaud's model.

Usually, proving an aging result consists in finding a two-point function $F(t_w, t_w + t)$, a quantity that measures the behaviour of the system at time $t + t_w$ after it has aged for the time t_w , such that a nontrivial limit

$$\lim_{\substack{t \to \infty \\ t/t_w = \theta}} F(t_w, t_w + t) = F(\theta)$$
(3)

exists. The choice of the two-point function is crucial. For instance it has be shown by Rinn, Maass, and Bouchaud (2000) that a good choice is

$$R(t_w, t_w + t) = \mathbb{EP}(X(t + t_w) = X(t_w)|E), \tag{4}$$

which is the probability that the system will be in the same state at the end of the observation period (i.e. at time $t + t_w$) as it was in the beginning (i.e. at time t_w). Another quantity exhibiting aging behaviour, which was studied by Fontes, Isopi, and Newman (2002) is

$$R^{q}(t_{w}, t_{w} + t) = \mathbb{E} \sum_{i \in \mathbb{Z}} [\mathbb{P}(X(t + t_{w}) = i | E, X(t_{w}))]^{2},$$
 (5)

which is the probability that two independent walkers will be at the same site after time $t + t_w$ if they were at the same site at time t_w , averaged over the distribution of the common starting point $X(t_w)$. These authors have proved that, for these two two-point functions, aging occurs when a = 0. We extend this result to the case a > 0. The

limiting object will be independent of a. Thus the parameter a could seem to be of no relevance for aging.

However, it is not the case for all two-point functions. For instance, for the function

$$\Pi(t_w, t_w + t) = \mathbb{EP}(X(t') = X(t_w) \forall t' \in [t_w, t_w + t] | E), \tag{6}$$

giving the probability that the system does not change its state between t_w and $t_w + t$, it was predicted by Rinn, Maass, and Bouchaud (2000) that there exists a constant γ such that the limit $\lim_{t_w \to \infty} \Pi(t_w, t_w + \theta t_w^{\gamma})$ exists and depends non-trivially on a. The name subaging was introduced for this type of behaviour, i.e. for the fact that there exists a constant $0 < \gamma < 1$ such that for some two-point function $F(t_w, t_w + t)$, there is a nontrivial limit

$$\lim_{\substack{t \to \infty \\ t \neq t_w^* = \theta}} F(t_w, t_w + t) = F(\theta). \tag{7}$$

One of the main results of the present paper is the proof of the subaging behaviour of the function (6) for an arbitrary $a \in [0, 1]$.

Let us have a closer look at the role of the parameter a. If a = 0, the dynamics of the model is sometimes referred as "Random hopping time (RHT) dynamics" (cf. Mathieu, 2000). In this case the rates w_{ij} do not depend on the value of E_j . Hence, the system jumps to all neighbouring sites with the same probability and the process X(t) can be regarded as a time change of the simple random walk.

On the other hand, if a > 0, the system is attracted to the deepest traps and the underlying discrete time Markov chain is some kind of random walk in a random environment (RWRE). There are already some results about aging of RWRE in dimension one (Dembo, Guionnet, and Zeitouni, 2001). It that article Sinai's RWRE is considered. It is proved there that there is aging on the scale $\log t/\log t_w \to \mathrm{const.}$

In our situation the energy landscape, far from being seen as a two-sided Brownian motion as in Sinai's RWRE, should be seen as essentially flat with few very narrow deep holes around the deep traps. The drifts on neighbouring sites are dependent and this dependency does not allow the existence of large domains with drift in one direction. This can be easily seen by looking at sites surrounding one particularly deep trap E_i . Here, the drift at site i-1 pushes the system very strongly to the right and at site i+1 to the left because the system is attracted to the site i. Moreover, these drifts have approximately the same size. A more precise description of this picture will be presented later (Section 5). However, these differences do not change notably the mechanism responsible for aging. Again, during the exploration of the random landscape, the process X finds deeper and deeper traps that slow down its dynamics.

It was observed numerically by Rinn, Maass, and Bouchaud (2000) that X(t) ages only if the temperature is low enough, $\beta > 1$. (In

the sequel we will consider only the low temperature regime.) This heuristically corresponds to the fact that if a=0 and $\beta>1$, the mean time $\mathbb{E}(\exp(\beta E_0))$ spent by X(t) at arbitrary site becomes infinite. This implies that the distribution of the depth at which we find the system at time t does not converge as $t \to \infty$. The process X(t) can find deeper and deeper traps where it stays longer.

If a > 0, the previous explanation is not precise. The time before the jump is shortened when a increases. On the other hand, the system is attracted to deep traps. This means that instead of staying in one deep trap, the process prefers to jump out and then to return back very quickly. For the two-point functions (4) and (5) these two effects cancel and the limiting behaviour is thus independent of a. For the two-point function (6), there cannot be cancellation, because the attraction to deeper traps has no influence on it.

Before stating the known results about the model we generalise it slightly. All statements in this paper do not actually require E_i to be an exponential random variable. The only property of E_i that we will need is that the random variable $\exp(\beta E_i)$ is in the domain of attraction of the totally asymmetric stable law with index $\beta^{-1} \equiv \alpha$. Clearly, the original exponential random variable satisfies this property.

Recently, this model was studied rigorously by Fontes, Isopi, and Newman (1999, 2002) in connection with the random voter model and chaotic time dependence. In this paper only the RHT case, a=0, was considered. If d=1 and $\beta>1$, they proved that the Markov chain X(t) possess an interesting property called there localisation. Namely, it was shown there that

$$\limsup_{t \to \infty} \sup_{i \in \mathbb{Z}} \mathbb{P}(X(t) = i|E) > 0 \qquad \mathbb{P}\text{-a.s.}$$
 (8)

Also aging for the two-point functions (4) and (5) was proved there. In dimension $d \geq 2$, results of this paper imply that there is no localisation in the sense of (8). However, there is numerical evidence (Rinn, Maass, and Bouchaud, 2000) that the system ages. A rigorous proof of this claim will be presented in a forthcoming paper (Ben Arous, Černý, and Mountford).

In this article we generalise the results of Fontes, Isopi, and Newman (2002) in dimension one to the general case, $a \neq 0$. As we have already noted, the main difficulty comes from the fact that the underlying discrete time Markov chain is not a simple random walk. We will prove aging for the quantities (4) and (5). We will then prove sub-aging for the two-point function (6).

As in Fontes, Isopi, and Newman (2002) we relate the asymptotic behaviour of quantities (4), (5), and (6) to the similar quantities computed using a singular diffusion Z(t) in a random environment ρ — singular meaning here that the single time distributions of Z are discrete.

Definition 1.1 (Diffusion with random speed measure). The random environment ρ is a random discrete measure, $\sum_i v_i \delta_{x_i}$, where the countable collection of (x_i, v_i) 's yields an inhomogeneous Poisson point process on $\mathbb{R} \times (0, \infty)$ with density measure $dx \alpha v^{-1-\alpha} dv$. Conditional on ρ , Z(s) is a diffusion process (with Z(0) = 0) that can be expressed as a time change of a standard one-dimensional Brownian motion W(t) with the speed measure ρ . Denoting $\ell(t, y)$ the local time of W(t) at y, we define

$$\phi^{\rho}(t) = \int \ell(t, y) \rho(dy) \tag{9}$$

and the stopping time $\psi^{\rho}(s)$ as the first time t when $\phi^{\rho}(t) = s$; then $Z(t) = W(\psi^{\rho}(t))$.

A more detailed description of time changes of Brownian motion can be found in Section 2.

Our main result about aging is the following

Theorem 1.2. For any $\beta > 1$ and $a \in [0,1]$ there exist nontrivial functions $R(\theta)$, $R^q(\theta)$ such that

$$\lim_{t \to \infty} R(t, t + \theta t) = \lim_{t \to \infty} \mathbb{EP}[X((1 + \theta)t) = X(t)|E] = R(\theta),$$

$$\lim_{t \to \infty} R^q(t, t + \theta t) = \lim_{t \to \infty} \mathbb{E} \sum_{i \in \mathbb{Z}} [\mathbb{P}(X((1 + \theta)t) = i|E, X(t))]^2 = R^q(\theta).$$
(10)

Moreover, $R(\theta)$ and $R^q(\theta)$ can be expressed using the similar quantities defined using the singular diffusion Z:

$$R(\theta) = \mathbb{EP}[Z(1+\theta) = Z(1)|\rho],$$

$$R^{q}(\theta) = \mathbb{E}\sum_{x \in \mathbb{P}} [\mathbb{P}(Z(1+\theta) = x|\rho, Z(1))]^{2}.$$
(11)

For a=0, this result is contained in Fontes, Isopi, and Newman (2002). Since the diffusion Z(t) does not depend on a, the functions $R(\theta)$ and $R^q(\theta)$ do not depend on it either. This is the result of the compensation of shorter visits of deep traps by the attraction to them.

We will also prove sub-aging for the quantity $\Pi(t_w, t_w + t)$. We use γ to denote the *subaging exponent*

$$\gamma = \frac{1}{1+\alpha} = \frac{\beta}{1+\beta}.\tag{12}$$

Theorem 1.3. For any $\beta > 1$ and $a \in [0,1]$ there exist a nontrivial function $\Pi(\theta)$ such that

$$\lim_{t \to \infty} \Pi(t, t + f_a(t, \theta)) = \lim_{t \to \infty} \mathbb{EP}[X(t') = X(t) \forall t' \in [t, t + f_a(t, \theta)] | E] = \Pi(\theta),$$
(13)

where the function f_a is given by

$$f_a(t,\theta) = \theta t^{\gamma(1-a)} L(t)^{1-a}, \tag{14}$$

and L(t) is a slowly varying function that is determined only by the distribution of E_0 . Its precise definition is given in Lemma 8.1. The function $\Pi(\theta)$ can be again written using the singular diffusion Z,

$$\Pi(\theta) = \int_0^\infty g_a^2(\theta u^{a-1}) dF(u), \tag{15}$$

where $F(u) = \mathbb{EP}[\rho(Z(1)) \leq u|\rho]$, and where $g_a(\lambda)$ is the Laplace transform $\mathbb{E}(e^{-\lambda T_a})$ of the random variable

$$T_a = 2^{a-1} \exp(a\beta E_0) [\mathbb{E}(\exp(-2a\beta E_0))]^{1-a}.$$
 (16)

If a = 0, (15) can be written as

$$\Pi(\theta) = \int_0^\infty e^{-\theta/u} dF(u). \tag{17}$$

Remark. Note that if E_i 's are exponential random variables, the function L(t) satisfies $L(t) \equiv 1$. The same is true if $\exp(\beta E_i)$ has a stable law.

As can be seen, in this case the function $\Pi(\theta)$ depends on a. This is not surprising since the compensation by attraction has no influence here and the jumps rates clearly depend on a.

This behaviour of the two-point functions $\Pi(t_w, t + t_w)$ and $R(t_w, t + t_w)$ is not difficult to understand, at least heuristically. One should first look at the behaviour of the distribution of the depth of the location of the process at time t_w . It can be proved that this depth grows like $t_w^{1/(1+\alpha)}$ (see Proposition 8.2). From this one can see that the main contribution to quantities (4) and (5) comes from trajectories of X(t) that, between times t_w and $t_w + t$, leave $t_w^{(a+\alpha)/(1+\alpha)}$ times the original site and then return to it. Each visit of the original site lasts an amount of time of order $t_w^{(1-a)/(1+\alpha)}$.

In the case of the two-point function (6), we are interested only in the first visit and thus the time t should scale as $t_w^{1/(1+\alpha)}$. Proofs can be found in Sections 7, 8 and 9. In Section 2 we summarise some known results about time-scale changes of Brownian motion and about point-process convergence. In Section 3 we express the process X and its scaled versions as a time-scale change and in Section 4 we introduce a coupling between the different scales of X. In Section 5 we prove convergence of speed measures which is used for time-scale change and we apply this result to show the convergence of finite time distributions of rescaled versions of X to the finite time distributions of Z.

2. Definitions and known results

In this section we define some notations that we will use often later, and we summarise some known results. 2.1. **Time-scale change of Brownian motion.** The limiting quantities $R(\theta)$, $R^q(\theta)$, and $\Pi(\theta)$ are expressed using the singular diffusion defined by a time change of Brownian motion. So, it will be convenient to express also the chains with discrete state space as a time-scale change of Brownian motion. The scale change is necessary if $a \neq 0$, because the process X(t) does not jump left or right with equal probabilities.

Consider a locally finite measure

$$\mu(dx) = \sum_{i} w_i \delta_{y_i}(dx) \tag{18}$$

which has atoms with weights w_i at positions y_i . The measure μ will be referred to as the speed measure. Let S be a strictly increasing function defined on the set $\{y_i\}$. We call such S the scaling function. Let us introduce slightly nonstandard notation $S \circ \mu$ for the "scaled measure"

$$(S \circ \mu)(dx) = \sum_{i} w_i \delta_{S(y_i)}(dx). \tag{19}$$

We use W(t) to denote the standard Brownian motion starting at 0. Let $\ell(t,y)$ be its local time. We define the function

$$\phi(\mu, S)(t) = \int_{\mathbb{R}} \ell(t, y)(S \circ \mu)(dy). \tag{20}$$

and the stopping time $\psi(\mu, S)(s)$ as the first time when $\phi(\mu, S)(t) = s$. The function $\phi(\mu, S)(t)$ is a nondecreasing, continuous function, and $\psi(\mu, S)(s)$ is its generalised right continuous inverse. It is an easy corollary of the results of Stone (1963) that the process

$$X(\mu, S)(t) = S^{-1}(W(\psi(\mu, S)(t)))$$
(21)

is a time-changed nearest neighbours random walk on the set of atoms of μ . Moreover, every nearest neighbours random walk on a countable, nowhere dense subset of $\mathbb R$ satisfying some mild conditions on transition probabilities can be expressed in this way. We call the process $X(\mu, S)$ the time-scale change of Brownian motion. If S is the identity function, we speak only about time change.

The following proposition describes the properties of $X(\mu, S)$ if the set of atoms of μ has no accumulation point. In this case we can suppose that the locations of atoms y_i satisfy $y_i < y_j$ if i < j. The claim is the consequence of Stone (1963, Section 3). The extra factor 2 comes from the fact that Stone uses the Brownian motion with generator $-\Delta$.

Proposition 2.1. The process $X(\mu, S)(t)$ is a nearest neighbours random walk on the set $\{y_i\}$ of atoms of μ . The waiting time in the state y_i is exponentially distributed with mean

$$2w_i \frac{(S(y_{i+1}) - S(y_i))(S(y_i) - S(y_{i-1}))}{S(y_{i+1}) - S(y_{i-1})}.$$
 (22)

After leaving state y_i , $X(\mu, S)$ enters states y_{i-1} and y_{i+1} with respective probabilities

$$\frac{S(y_{i+1}) - S(y_i)}{S(y_{i+1}) - S(y_{i-1})} \quad and \quad \frac{S(y_i) - S(y_{i-1})}{S(y_{i+1}) - S(y_{i-1})}.$$
 (23)

It will be useful to introduce another process $Y(\mu, S)$ as

$$Y(\mu, S)(t) = X(S \circ \mu, \mathrm{Id})(t), \tag{24}$$

where Id is the identity function on \mathbb{R} . The process $Y(\mu, S)$ can be regarded as $X(\mu, S)$ before the final change of scale in (21). Actually,

$$Y(\mu, S)(t) = W(\psi(\mu, S)(t)). \tag{25}$$

We will also need processes that are not started at the origin but at some point $x \in \text{supp } \mu$. They are defined in the obvious way using the Brownian motion started at S(x). We use $X(\mu, S; x)$ and $Y(\mu, S; x)$ to denote them.

2.2. Point process convergence. To be able to work with quantities (4)–(6) that have a discrete nature (in the sense that they depend on the probability being exactly at some place) we recall the definition of the point process convergence of measures introduced in Fontes, Isopi, and Newman (2002). Let \mathcal{M} denote the set of locally finite measures on \mathbb{R} .

Definition 2.2 (Fontes, Isopi, and Newman (2002)). Given a family $\nu, \nu^{\varepsilon}, \varepsilon > 0$, in \mathcal{M} , we say that ν^{ε} converges in the point process sense to ν , and write $\nu^{\varepsilon} \stackrel{pp}{\to} \nu$, as $\varepsilon \to 0$, provided the following holds: if the atoms of ν , ν^{ε} are, respectively, at the distinct locations $y_i, y_{i'}^{\varepsilon}$ with weights $w_i, w_{i'}^{\varepsilon}$, then the subsets of $V^{\varepsilon} \equiv \cup_{i'} \{(y_{i}^{\varepsilon}, w_{i'}^{\varepsilon})\}$ of $\mathbb{R} \times (0, \infty)$ converge to $V \equiv \cup_i \{(y_i, w_i)\}$ as $\varepsilon \to 0$ in the sense that for any open U, whose closure \bar{U} is a compact subset of $\mathbb{R} \times (0, \infty)$ such that its boundary contains no points of V, the number of points $|V^{\varepsilon} \cap U|$ in $V^{\varepsilon} \cap U$ is finite and equals $|V \cap U|$ for all ε small enough.

Beside this type of convergence we will use the following two more common types of convergence

Definition 2.3. For the same family as in the previous definition, we say that ν^{ε} converges vaguely to ν , and write $\nu^{\varepsilon} \stackrel{v}{\to} \nu$, as $\varepsilon \to 0$, if for all continuous real-valued functions f on \mathbb{R} with bounded support $\int f(y)\nu^{\varepsilon}(dy) \to \int f(y)\nu(dy)$ as $\varepsilon \to 0$. We say that ν^{ε} converges weakly, and we write $\nu^{\varepsilon} \stackrel{w}{\to} \nu$, as $\varepsilon \to 0$, if the same is true for all bounded continuous functions on \mathbb{R} .

To prove the point process convergence we will use the next lemma that is the copy of Proposition 2.1 of Fontes, Isopi, and Newman (2002).

Let ν , ν^{ε} be locally finite measures on \mathbb{R} and let (y_i, w_i) , $(y_i^{\varepsilon}, w_i^{\varepsilon})$ be the sets of atoms of these measures (y_i) is the position and w_i is the weight of the atom).

Condition 1. For each l there exists a sequence $j_l(\varepsilon)$ such that

$$(y_{j_l(\varepsilon)}^{\varepsilon}, w_{j_l(\varepsilon)}^{\varepsilon}) \to (y_l, w_l) \quad as \quad \varepsilon \to 0.$$
 (26)

- **Lemma 2.4.** For any family ν , ν^{ε} , $\varepsilon > 0$, in \mathcal{M} , the following two assertions hold. If $\nu^{\varepsilon} \xrightarrow{pp} \nu$ as $\varepsilon \to 0$, then Condition 1 holds. If Condition 1 holds and $\nu^{\varepsilon} \xrightarrow{v} \nu$ as $\varepsilon \to 0$, then also $\nu^{\varepsilon} \xrightarrow{pp} \nu$ as $\varepsilon \to 0$.
- 2.3. Convergence of the fixed time distributions. We want to formulate, for the future use, a series of results of Fontes, Isopi, and Newman (2002). They will allow us to deduce the convergence of fixed time distributions from the convergence of speed measures.

Proposition 2.5. Let μ^{ε} , μ be the collection of deterministic locally finite measures, and let Y^{ε} , Y be defined by

$$Y^{\varepsilon}(t) = Y(\mu^{\varepsilon}, Id)(t) \quad and \quad Y(t) = Y(\mu, Id)(t).$$
 (27)

For any deterministic $t_0 > 0$, let ν^{ε} denote the distribution of $Y^{\varepsilon}(t_0)$ and ν denote the distribution of $Y(t_0)$. Suppose

$$\mu^{\varepsilon} \xrightarrow{v} \mu \quad and \quad \mu^{\varepsilon} \xrightarrow{pp} \mu \quad as \quad \varepsilon \to 0.$$
 (28)

(i) Then, as $\varepsilon \to 0$,

$$\nu^{\varepsilon} \xrightarrow{v} \nu \quad and \quad \nu^{\varepsilon} \xrightarrow{pp} \nu.$$
 (29)

(ii) Let $(x_i^{\varepsilon}, v_i^{\varepsilon})$ and (x_i, v_i) be the collections of atoms of μ^{ε} and μ . Similarly, let $(y_i^{\varepsilon}, w_i^{\varepsilon})$ and (y_i, w_i) be the collections of atoms of ν^{ε} and ν . Then the sets of locations of the atoms are equal,

$$\{y_i^{\varepsilon}\} = \{x_i^{\varepsilon}\} \quad and \quad \{y_i\} = \{x_i\}.$$
 (30)

(iii) Suppose that we have denoted x_i 's and y_i 's in such a way that $x_i = y_i$, $x_i^{\varepsilon} = y_i^{\varepsilon}$ (which is possible by (ii)). Let the sequence $j_l(\varepsilon)$ satisfy

$$(x_{j_l(\varepsilon)}^{\varepsilon}, v_{j_l(\varepsilon)}^{\varepsilon}) \to (x_l, v_l) \quad as \quad \varepsilon \to 0.$$
 (31)

Then the sequence of corresponding atoms of ν^{ε} satisfies

$$(y_{j_l(\varepsilon)}^{\varepsilon}, w_{j_l(\varepsilon)}^{\varepsilon}) = (x_{j_l(\varepsilon)}^{\varepsilon}, w_{j_l(\varepsilon)}^{\varepsilon}) \to (y_l, w_l) \quad as \quad \varepsilon \to 0.$$
 (32)

- (iv) Let $z^{\varepsilon} \to z$ and $t^{\varepsilon} \to t_0$ as $\varepsilon \to 0$. Then parts (i)-(iii) stay valid if we replace the process $Y^{\varepsilon}(t)$ by the process started outside the origin $Y(\mu^{\varepsilon}, Id; z^{\varepsilon})$, the process Y(t) by $Y(\mu, Id; z)$, and we define ν^{ε} as the distribution of $Y^{\varepsilon}(t^{\varepsilon})$.
- Part (i) of this proposition is stated as Theorem 2.1 in Fontes, Isopi, and Newman (2002). Part (ii) is a consequence of Lemmas 2.1 and 2.3 of the same paper. Part (iii) follows from the proof of that theorem, but it is not stated there explicitly. Its proof is, however, the central part of the proof of (i). The remaining part is an easy consequence of (i)–(iii) and of the joint continuity of the local time $\ell(t, y)$.

3. Expression of X(t) in terms of Brownian motion

To explore the asymptotic behaviour of the chain X(t), we consider its scaling limit

$$X^{\varepsilon}(t) = \varepsilon X(t/\varepsilon c_{\varepsilon}). \tag{33}$$

The constant c_{ε} will be determined later. For the time being the reader can consider $c_{\varepsilon} \sim \varepsilon^{1/\alpha}$.

As we already noted in the previous section, it is convenient to express the walks X(t) and $X^{\varepsilon}(t)$ as a time-scale change of the standard Brownian motion W(t) started at 0. To achieve it we use Proposition 2.1. We define measures

$$\mu(dx) = \mu^{1}(dx) = \sum_{i \in \mathbb{Z}} \tau_{i} \delta_{i}(dx) \quad \text{and} \quad \mu^{\varepsilon}(dx) = c_{\varepsilon} \sum_{i \in \mathbb{Z}} \tau_{i} \delta_{\varepsilon i}(dx), \quad (34)$$

where

$$\tau_i = \frac{1}{2} \exp(\beta E_i) \mathbb{E}(\exp(-2a\beta E_0)). \tag{35}$$

We will consider the following scaling function. Let

$$r_i = \frac{\exp(-\beta a(E_i + E_{i+i}))}{\mathbb{E}(\exp(-2\beta a E_0))},$$
(36)

and let

$$S(i) = \begin{cases} \sum_{j=0}^{i-1} r_j & \text{if } i \ge 0, \\ -\sum_{j=i}^{-1} r_j & \text{otherwise.} \end{cases}$$
 (37)

The constant factor $\mathbb{E}(\exp(-2\beta aE_0))$ that appears in (35) and (36) is not substantial, but it is convenient and it will simplify some expressions later.

We use $\tilde{X}^{\varepsilon}(t)$, $0 < \varepsilon \le 1$, to denote the process

$$\tilde{X}^{\varepsilon}(t) = X(\mu^{\varepsilon}, \varepsilon S(\varepsilon^{-1}\cdot))(t),$$
 (38)

which means that $\tilde{X}(t)$ is time-scale change of Brownian motion with speed measure μ^{ε} and scale function $\varepsilon S(\varepsilon^{-1}\cdot)$. If we write $\psi^{\varepsilon}(t)$ for $\psi(\mu^{\varepsilon}, \varepsilon S(\varepsilon^{-1}\cdot))(t)$, we have

$$\tilde{X}^{\varepsilon}(t) = \varepsilon S^{-1}(\varepsilon^{-1}W^{\varepsilon}(\psi^{\varepsilon}(t))). \tag{39}$$

The process W^{ε} is the rescaled Brownian motion, $W^{\varepsilon}(t) = \varepsilon W(\varepsilon^{-2}t)$, which has the same distribution as W(t). It is introduced only to simplify the proof of the next lemma. In the sequel we will omit the superscript if $\varepsilon = 1$, i.e. we will write $\tilde{X}(t)$ for $\tilde{X}^{1}(t)$, etc. Note that the function $S^{-1}(\cdot)$ is well defined for all values of its argument. Indeed, the set of atoms of $\varepsilon S(\varepsilon^{-1}\cdot) \circ \mu^{\varepsilon}$ is the set $\{\varepsilon S(i) : i \in \mathbb{Z}\}$, and thus $\varepsilon^{-1}W^{\varepsilon}(\psi^{\varepsilon}(t))$ takes values only in $\{S(i) : i \in \mathbb{Z}\}$.

Proposition 3.1. The processes $\tilde{X}(t)$ and $\tilde{X}^{\varepsilon}(t)$ have the same distribution as X(t) and $X^{\varepsilon}(t) = \varepsilon X(t/c_{\varepsilon}\varepsilon)$.

Proof. We use the symbol \sim to denote the equality in distribution. The time that X(t) stays at site i is exponentially distributed with mean $(w_{i,i+1} + w_{i,i-1})^{-1}$. The probability that it jumps right or left is

$$\frac{w_{i,i+1}}{w_{i,i+1} + w_{i,i-1}}$$
 and $\frac{w_{i,i-1}}{w_{i,i+1} + w_{i,i-1}}$. (40)

Plugging the definition (2) of w_{ij} into these expressions, it is easy to see that these values coincide with the same quantities for $\tilde{X}(t)$ which can be computed using Proposition 2.1. This implies that $X(t) \sim \tilde{X}(t)$.

To compare the distributions of $X^{\varepsilon}(t)$ and $X^{\varepsilon}(t)$, let us first look at the scaling of $\psi^{\varepsilon}(t)$. After an easy calculation, using the fact that the local time $\ell^{\varepsilon}(t,y)$ of W^{ε} satisfies $\ell^{\varepsilon}(t,y) = \varepsilon \ell(\varepsilon^{-2}t, \varepsilon^{-1}y)$, we obtain

$$\phi^{\varepsilon}(t) = \int \ell^{\varepsilon}(t, y) (\varepsilon S(\varepsilon^{-1} \cdot) \circ \mu^{\varepsilon}) (dy) = \varepsilon c_{\varepsilon} \phi(\varepsilon^{-2} t). \tag{41}$$

From it we get $\psi^{\varepsilon}(t) = \varepsilon^2 \psi(t/\varepsilon c_{\varepsilon})$. Hence,

$$\varepsilon \tilde{X}(t/\varepsilon c_{\varepsilon}) = \varepsilon S^{-1}(W(\psi(t/\varepsilon c_{\varepsilon}))) = \varepsilon S^{-1}(W(\varepsilon^{-2}\psi^{\varepsilon}(t))) =$$

$$= \varepsilon S^{-1}(\varepsilon^{-1}W^{\varepsilon}(\psi^{\varepsilon}(t))) = \tilde{X}^{\varepsilon}(t),$$
(42)

where we used the scaling of W(t) and (39). Since $\tilde{X}(t)$ has the same distribution as X(t), the same is valid for $\tilde{X}^{\varepsilon}(t)$ and $X^{\varepsilon}(t)$.

4. A COUPLING FOR WALKS ON DIFFERENT SCALES

It is convenient to introduce the processes Y(t) and $Y^{\varepsilon}(t)$ that are only a time change of Brownian motion with speed measures $S \circ \mu$ and $\varepsilon S(\varepsilon^{-1} \cdot) \circ \mu^{\varepsilon}$. Namely,

$$Y^{\varepsilon}(t) = Y(\mu^{\varepsilon}, \varepsilon S(\varepsilon^{-1} \cdot))(t)$$
 and $Y(t) = Y(\mu, S)(t)$. (43)

Using (25) we have

$$Y(t) = W(\psi(t))$$
 and $Y^{\varepsilon}(t) = W(\psi^{\varepsilon}(t)).$ (44)

The original processes X and X^{ε} are related to them by

$$X(t) = S^{-1}(Y(t))$$
 and $X^{\varepsilon}(t) = \varepsilon S^{-1}(\varepsilon^{-1}Y^{\varepsilon}(t)).$ (45)

In the sequel we want to use Proposition 2.5 to prove the convergence of the finite time distributions of Y^{ε} . Thus, we want to apply this proposition to the sequence of random speed measures μ^{ε} . It is easy to see that convergence in distribution of this sequence is not sufficient for its application. That is why we will construct a coupling between measures μ^{ε} on different scales ε on a larger probability space. Using this coupling we obtain the a.s. convergence on this space. It is not surprising that the same coupling as in Fontes, Isopi, and Newman (2002) does the job.

Consider the Lévy process $V(x), x \in \mathbb{R}, V(0) = 0$, with stationary and independent increments and cadlag paths defined on $(\bar{\Omega}, \bar{\mathcal{F}}, \bar{\mathbb{P}})$ given by

$$\bar{\mathbb{E}}\left[e^{ir(V(x+x_0)-V(x_0))}\right] = \exp\left[x\alpha \int_0^\infty (e^{irw} - 1)w^{-1-\alpha}dw\right]. \tag{46}$$

Let $\bar{\rho}$ be the random Lebesgue-Stieltjes measure on \mathbb{R} associated to V, i.e. $\bar{\rho}(a,b] = V(b) - V(a)$. It is a known fact that $\bar{\rho}(dx) = \sum_j v_j \delta_{x_j}(dx)$, where (x_j,v_j) is an inhomogeneous Poison point process with density $dx \, \alpha v^{-1-\alpha} dv$. Note that $\bar{\rho}$ has the same distribution as ρ which we used as speed measure in the definition of the singular diffusion Z.

For each fixed $\varepsilon > 0$, we will now define the sequence of i.i.d. random variables E_i^{ε} such that E_i^{ε} 's are defined on the same space as V and $\bar{\rho}$ and they have the same distribution as E_0 .

Define a function $G:[0,\infty)\mapsto [0,\infty)$ such that

$$\bar{\mathbb{P}}(V(1) > G(x)) = \mathbb{P}(\tau_0 > x). \tag{47}$$

The function G is well-defined since V(1) has continuous distribution, it is nondecreasing and right continuous, and hence has nondecreasing right-continuous generalised inverse G^{-1} . Let $g_{\varepsilon}:[0,\infty)\mapsto [0,\infty)$ be defined as

$$g_{\varepsilon}(x) = c_{\varepsilon} G^{-1}(\varepsilon^{-1/\alpha} x) \quad \text{for all} \quad x \ge 0,$$
 (48)

where

$$c_{\varepsilon} = \left(\inf[t \ge 0 : \mathbb{P}(\tau_0 > t) \le \varepsilon]\right)^{-1}.$$
 (49)

Note that if τ_0 is the α stable random variable with characteristic function

$$\mathbb{E}(e^{ir\tau_0}) = \exp\left[\alpha \int_0^\infty (e^{irw} - 1)w^{-1-\alpha}dw\right],\tag{50}$$

the choice of c_{ε} and g_{ε} can be simplified (although it does not correspond to the previous definition)

$$c_{\varepsilon} = \varepsilon^{1/\alpha}$$
 and $g_{\varepsilon}(y) \equiv y$. (51)

The reader who is not interested in the technical details should keep this choice in mind.

Lemma 4.1. Let

$$\tau_i^{\varepsilon} = \frac{1}{c_{\varepsilon}} g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)), \tag{52}$$

and

$$E_i^{\varepsilon} = \frac{1}{\beta} \log \left(\frac{2\tau_i^{\varepsilon}}{\mathbb{E}(\exp(-2a\beta E_0))} \right). \tag{53}$$

Then for any $\varepsilon > 0$, the τ_i^{ε} are i.i.d. with the same law as τ_0 , and $\{E_i^{\varepsilon}\}_{i \in \mathbb{Z}}$ have the same distribution as $\{E_i\}_{i \in \mathbb{Z}}$.

Proof. By stationarity and independence of increments of V it is sufficient to show $\bar{\mathbb{P}}(\tau_0^{\varepsilon} > t) = \mathbb{P}(\tau_0 > t)$. However,

$$\bar{\mathbb{P}}(\tau_0^{\varepsilon} > t) = \bar{\mathbb{P}}(V(\varepsilon) > \varepsilon^{1/\alpha} G(t)) \tag{54}$$

by the definitions of τ_0^{ε} and G. The result then follows from (47) and the scaling invariance of $V: V(\varepsilon) \sim \varepsilon^{1/\alpha} V(1)$. The second claim follows easily using (35).

Let us now define the random speed measures $\bar{\mu}^{\varepsilon}$ using the collections $\{E_i^{\varepsilon}\}$ from the previous lemma,

$$\bar{\mu}^{\varepsilon}(dx) = \sum_{i \in \mathbb{Z}} c_{\varepsilon} \tau_i^{\varepsilon} \delta_{\varepsilon i}(dx). \tag{55}$$

We also define the scaling functions S_{ε} similarly as in (37). Let

$$r_i^{\varepsilon} = \frac{\exp(-\beta a(E_i^{\varepsilon} + E_{i+1}^{\varepsilon}))}{\mathbb{E}(\exp(-2a\beta E_0))}$$
 (56)

and

$$S_{\varepsilon}(i) = \begin{cases} \sum_{j=0}^{i-1} r_j^{\varepsilon} & \text{if } i \ge 0, \\ -\sum_{j=i}^{-1} r_j^{\varepsilon} & \text{otherwise.} \end{cases}$$
 (57)

It is an easy consequence of Lemma 4.1 that $\bar{\mu}^{\varepsilon} \sim \mu^{\varepsilon}$ and $S_{\varepsilon} \sim S$ for any $\varepsilon \in (0,1]$.

5. Convergence of speed measures

The following proposition proves the convergence of the scaled speed measures. If S is the identity, i.e. a=0, it corresponds to Proposition 3.1 of Fontes, Isopi, and Newman (2002).

Proposition 5.1. Let $\bar{\mu}^{\varepsilon}$ and $\bar{\rho}$ be defined as above. Then

$$\varepsilon S_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \bar{\mu}^{\varepsilon} \xrightarrow{v} \bar{\rho} \quad and \quad \varepsilon S_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \bar{\mu}^{\varepsilon} \xrightarrow{pp} \bar{\rho} \quad as \ \varepsilon \to 0 \quad \bar{\mathbb{P}}\text{-}a.s. (58)$$

The proof requires three technical lemmas.

Lemma 5.2. As $\varepsilon \to 0$ we have

$$\varepsilon S_{\varepsilon}(\lfloor \varepsilon^{-1} y \rfloor) \to y \quad as \ \varepsilon \to 0 \quad \bar{\mathbb{P}} \text{-}a.s.$$
 (59)

uniformly on compact intervals.

Notice that this lemma sheds more light on the difference between the discrete time embedded walk of the process X and the Sinai's RWRE. In the case of Sinai's RWRE the scale function S corresponds, loosely speaking, to the function

$$S'(n) = \sum_{i=1}^{n} \rho_1 \dots \rho_n, \tag{60}$$

where $\rho_i = (1 - p_i)/p_i$, p_i is the probability going right at i, and p_i 's are i.i.d. In our case $\rho_i = r_i/r_{i-1}$. An easy computation gives that

the product $\rho_1 \dots \rho_n$ depends only on E_0 and E_{n+1} . Thus, S'(n) is in our situation essentially a sum of i.i.d. random variables which is definitively not the case for the Sinai's RWRE.

Proof of Lemma 5.2. We consider only y > 0. The proof for y < 0 is very similar. By definition of S_{ε} we have $\varepsilon S_{\varepsilon}(\lfloor \varepsilon^{-1}y \rfloor) = \varepsilon \sum_{j=0}^{\lfloor \varepsilon^{-1}y \rfloor - 1} r_j^{\varepsilon}$, where for fixed ε the sequence r_i^{ε} is an ergodic sequence of bounded positive random variables. Moreover, r_i^{ε} is independent of all r_j^{ε} with $j \notin \{i-1,i,i+1\}$. The $\overline{\mathbb{P}}$ -a.s. convergence for fixed y is then a consequence of the strong law of large numbers for triangular arrays. Note that this law of large numbers can be easily proved in our context using the standard methods, because the variables r_i^{ε} are bounded and thus their moments of arbitrary large degree are finite. The uniform convergence on compact intervals is easy to prove using the fact that $S_{\varepsilon}(i)$ is increasing and the identity function is continuous.

The next two lemmas correspond to Lemmas 3.1 and 3.2 of Fontes, Isopi, and Newman (2002). We state them without proofs.

Lemma 5.3. For any fixed y > 0, $g^{\varepsilon}(y) \to y$ as $\varepsilon \to 0$.

Lemma 5.4. For any $\delta' > 0$, there exist constants C' and C'' in $(0, \infty)$ such that

$$g_{\varepsilon}(x) \le C' x^{1-\delta'} \quad for \quad \varepsilon^{1/\alpha} \le x \le 1 \quad and \quad \varepsilon \le C''.$$
 (61)

Proof of Proposition 5.1. We first prove the vague convergence. Let f be a bounded continuous function with compact support $I \subset \mathbb{R}$. Then,

$$\int f(x)(\varepsilon S_{\varepsilon}(\varepsilon^{-1}x) \circ \bar{\mu}^{\varepsilon})(dx) = \sum_{i \in J_0^{\varepsilon}} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)),$$
(62)

where we used the notation

$$J_y^{\varepsilon} = \{ i \in \mathbb{Z} : \varepsilon S_{\varepsilon}(i) \in I, V(\varepsilon(i+1)) - V(\varepsilon i) \ge y \}.$$
 (63)

Choose now $\delta>0$. To estimate the last sum, we treat separately the sums over J^{ε}_{δ} , $J^{\varepsilon}_{\varepsilon^{1/\alpha}}\setminus J^{\varepsilon}_{\delta}$ and $J^{\varepsilon}_{0}\setminus J^{\varepsilon}_{\varepsilon^{1/\alpha}}$.

Due to the convergence of $\varepsilon S_{\varepsilon}(\varepsilon^{-1}\cdot)$ to the identity, we know that

Due to the convergence of $\varepsilon S_{\varepsilon}(\varepsilon^{-1}\cdot)$ to the identity, we know that for ε small enough there is a small neighbourhood I' of I such that $J_0^{\varepsilon} \subset \varepsilon^{-1}I'$. The process V has $\bar{\mathbb{P}}$ -a.s. only finitely many jumps larger than δ in I', so the first sum has only a finite number of terms. Using the continuity of f and applying Lemmas 5.2 and 5.3 we have

$$\sum_{i \in J_{\delta}^{\varepsilon}} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)) \to \sum_{j: v_j \ge \delta} f(x_j) v_j, \qquad (64)$$

with (x_i, v_i) being the set of atoms of $\bar{\rho}$. In the previous expression we also use the fact that $i\varepsilon \to x_i$ for the corresponding terms in the sums.

By Lemma 5.4 we have for some δ' such that $\delta' + \alpha \leq 1$

$$\sum_{i \in J_{\varepsilon^{1/\alpha}}^{\varepsilon} \setminus J_{\delta}^{\varepsilon}} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)) \leq \\
\leq C \sum_{i \in J_{\varepsilon^{1/\alpha}}^{\varepsilon} \setminus J_{\delta}^{\varepsilon}} \left(V(\varepsilon(i+1)) - V(\varepsilon i)\right)^{1-\delta'} \leq C \sum_{\substack{j: v_{j} \leq \delta \\ x_{j} \in I'}} v_{j}^{1-\delta'} = H_{\delta}.$$
(65)

From the definition of the point process (x_i, v_i) we have

$$\bar{\mathbb{E}}(H_{\delta}) \le \alpha |I'| \int_0^{\delta} w^{1-\delta'} w^{-1-\alpha} dw \to 0 \quad \text{as} \quad \delta \to 0.$$
 (66)

Since H_{δ} is decreasing and positive, the limit $\lim_{\delta\to 0} H_{\delta}$ exists $\bar{\mathbb{P}}$ -a.s. The dominated convergence theorem then gives $\bar{\mathbb{E}} \lim_{\delta\to 0} H_{\delta} = 0$, and thus $\lim_{\delta\to 0} H_{\delta} = 0$ $\bar{\mathbb{P}}$ -a.s.

The third part of the sum is also negligible for ε small enough. Indeed, by monotonicity of g_{ε} , we have $g_{\varepsilon}(x) \leq g_{\varepsilon}(\varepsilon^{1/\alpha}) \leq Cc_{\varepsilon}$ for all $x \leq \varepsilon^{1/\alpha}$. Hence,

$$\sum_{i \in J_0^{\varepsilon} \setminus J_{\varepsilon^{1/\alpha}}^{\varepsilon}} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)) \leq
\leq C' c_{\varepsilon} \sum_{i \in \varepsilon^{-1} I' \cap \mathbb{Z}} 1 \leq C'' c_{\varepsilon} \varepsilon^{-1} \to 0 \quad \text{as} \quad \varepsilon \to 0.$$
(67)

In the last equation we use the fact that if τ_0 is in the domain of attraction of the stable law with index α , there exists $\kappa > 0$ such that the function c_{ε} can be bounded from above by $C\varepsilon^{-\kappa+1/\alpha}$ with $-\kappa + 1/\alpha > 1$.

Putting now all three parts together, we have

$$\lim_{\varepsilon \to 0} \sum_{i \in J_0^{\varepsilon}} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon}(V(\varepsilon(i+1)) - V(\varepsilon i)) =$$

$$= \lim_{\delta \to 0} \sum_{j: v_{\delta} > \delta} f(x_j) v_j = \int f d\bar{\rho}.$$
(68)

This proves the vague convergence.

To prove the point process convergence we use Lemma 2.4. Since we have already proved the vague convergence, we must only verify Condition 1 for the measures $\varepsilon S_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \bar{\mu}^{\varepsilon}$ and $\bar{\rho}$. Thus, for any atom (x_l, v_l) of $\bar{\rho}$ we want to find a sequence $j_l(\varepsilon)$ such that

$$\varepsilon S_{\varepsilon}(j_l(\varepsilon)) \to x_l \quad \text{and} \quad g_{\varepsilon}(V(\varepsilon(j_l(\varepsilon)+1)) - V(\varepsilon j_l(\varepsilon))) \to v_l. \quad (69)$$

Choose $j_l(\varepsilon)$ such that $x_l \in (\varepsilon j_l(\varepsilon), \varepsilon (j_l(\varepsilon) + 1)]$. Then by Lemma 5.2 we have the first statement of (69), and by Lemma 5.3 we have the second. This finishes the proof of Proposition 5.1.

6. Change of scale for fixed time distributions

Write \bar{X}^{ε} and \bar{X} for the processes defined as in (38), but using the speed measures $\bar{\mu}^{\varepsilon}$ and the scaling functions S_{ε} . Since $\bar{\mu}^{\varepsilon} \sim \mu^{\varepsilon}$ and $S_{\varepsilon} \sim S$, we have $\bar{X}^{\varepsilon} \sim X^{\varepsilon}$. Similarly, we define the processes \bar{Y}^{ε} , \bar{Y} as in (44), and \bar{Z} as in Definition 1.1 using the measures with bars. Evidently, $\bar{Y}^{\varepsilon} \sim Y^{\varepsilon}$, $\bar{Y} \sim Y$ and $\bar{Z} \sim Z$. The following proposition is a consequence of Propositions 2.5 and 5.1.

Proposition 6.1. Fix $t_0 > 0$. Write $\bar{\nu}_{Y,V}^{\varepsilon}$ for the distribution of $\bar{Y}^{\varepsilon}(t_0)$ and $\bar{\nu}_V$ for the distribution of $\bar{Z}(t_0)$ conditionally on V. Then, $\bar{\mathbb{P}}$ -a.s we have

$$\bar{\nu}_{YV}^{\varepsilon} \xrightarrow{v} \bar{\nu}_{V} \quad and \quad \bar{\nu}_{YV}^{\varepsilon} \xrightarrow{pp} \bar{\nu}_{V} \quad as \quad \varepsilon \to 0.$$
 (70)

The proof of the convergence of the fixed time distribution of \bar{X}^{ε} will be finished if we can compare the limits of \bar{X}^{ε} and \bar{Y}^{ε} .

Proposition 6.2. Fix t_0 as in Proposition 6.1. Let $\bar{\nu}_{X,V}^{\varepsilon}$ denote the distribution of $\bar{X}^{\varepsilon}(t_0)$ conditionally on V. Then, $\bar{\mathbb{P}}$ -a.s. we have

$$\lim_{\varepsilon \to 0} \bar{\nu}_{X,V}^{\varepsilon} = \lim_{\varepsilon \to 0} \bar{\nu}_{Y,V}^{\varepsilon} = \bar{\nu}_{V}, \tag{71}$$

where the limits are taken in both the vague and the point process sense.

Proof. As an easy consequence of Lemma 5.2 we have

$$\varepsilon S_{\varepsilon}^{-1}(\varepsilon^{-1}y) \to y \qquad \bar{\mathbb{P}}\text{-a.s.}$$
 (72)

We will again apply Lemma 2.4 to prove the convergence. Let f be a continuous function with bounded support $I \subset \mathbb{R}$. By continuity of f and (72), choosing the fixed realisation of Brownian motion W, we have $\bar{\mathbb{P}}$ -a.s.

$$\lim_{\varepsilon \to 0} f(\bar{X}^{\varepsilon}(t_0)) = \lim_{\varepsilon \to 0} f(\bar{Y}^{\varepsilon}(t_0)). \tag{73}$$

A standard application of the dominated convergence theorem yields

$$\lim_{\varepsilon \to 0} \int f \, d\bar{\nu}_{X,V}^{\varepsilon} = \lim_{\varepsilon \to 0} \int f \, d\bar{\nu}_{Y,V}^{\varepsilon} = \int f \, d\bar{\nu}_{V}. \tag{74}$$

We finally verify Condition 1. Write $(x_i^{\varepsilon}, v_i^{\varepsilon})$, $(y_i^{\varepsilon}, w_i^{\varepsilon})$ for the collections of atoms of $\bar{\nu}_{X,V}^{\varepsilon}$ and $\bar{\nu}_{Y,V}^{\varepsilon}$. By Proposition 2.5(ii) we can choose $x_i^{\varepsilon} = \varepsilon i$ and $y_i^{\varepsilon} = \varepsilon S_{\varepsilon}(i)$, setting eventually v_i^{ε} , resp. w_i^{ε} , equal to zero if there is no atom at x_i^{ε} , resp. y_i^{ε} . Using this choice of x_i^{ε} and y_i^{ε} and the relation (45) we have $v_i^{\varepsilon} = w_i^{\varepsilon}$. Let (z_l, u_l) be the collection of atoms of $\bar{\nu}_V$ and $j_l(\varepsilon)$ be the sequence of indexes such that $(y_{j_l(\varepsilon)}, w_{j_l(\varepsilon)}) \to (z_l, u_l)$. Then by (72) we have $(x_{j_l(\varepsilon)}, v_{j_l(\varepsilon)}) \to (z_l, u_l)$ which completes the proof.

7. Proof of Theorem 1.2

We first express the quantities that we are interested in using the processes \bar{X}^{ε} . From the definition of \tilde{X}^{ε} , Proposition 3.1, and the fact that $\bar{X}^{\varepsilon} \sim \tilde{X}^{\varepsilon}$ we get that if the following limits exist (as we show below), they should satisfy

$$\lim_{t_w \to \infty} \mathbb{EP}\big[X((1+\theta)t_w) = X(t_w)|E\big] = \lim_{\varepsilon \to 0} \bar{\mathbb{E}}\bar{\mathbb{P}}\big[\bar{X}^{\varepsilon}(1+\theta) = \bar{X}^{\varepsilon}(1)|V\big] \equiv \lim_{\varepsilon \to 0} R_{\varepsilon}(\theta) \quad (75)$$

and similarly

$$\lim_{t_w \to \infty} \mathbb{E} \sum_{i \in \mathbb{Z}} \left[\mathbb{P}(X((1+\theta)t_w) = i|E, X(t_w)) \right]^2$$

$$\lim_{\varepsilon \to 0} \bar{\mathbb{E}} \sum_{i \in \mathbb{Z}} \left[\bar{\mathbb{P}}(\bar{X}^{\varepsilon}(1+\theta) = i\varepsilon|V, \bar{X}^{\varepsilon}(1)) \right]^2 \equiv \lim_{\varepsilon \to 0} R_{\varepsilon}^q(\theta). \quad (76)$$

We introduce some notation for the sets of atoms of the measures we will consider. In the following everything depends on the realisation of the Lévy process V and we will not denote this dependence explicitly. We write

$$\bar{\mu}^{\varepsilon} = \sum_{i} v_{i}^{\varepsilon} \delta_{x_{i}^{\varepsilon}} \quad \text{and} \quad \bar{\rho} = \sum_{i} v_{i} \delta_{x_{i}}.$$
(77)

The atoms of the distribution ν_1^{ε} of $\bar{X}^{\varepsilon}(1)$ will be denoted by $(x_i^{\varepsilon}, w_i^{\varepsilon})$. Similarly, (x_i, w_i) denotes the atoms of the distribution ν_1 of $\bar{Z}(1)$. The weights of the joint distribution of $\bar{X}^{\varepsilon}(1)$ and $\bar{X}^{\varepsilon}(1+\theta)$ will be denoted by w_{ii}^{ε} ,

$$w_{ij}^{\varepsilon} = \bar{\mathbb{P}}\big[(\bar{X}^{\varepsilon}(1) = x_i^{\varepsilon}) \cap (\bar{X}^{\varepsilon}(1+\theta) = x_j^{\varepsilon})|V\big],$$

$$w_{ij} = \bar{\mathbb{P}}\big[(\bar{Z}(1) = x_i) \cap (\bar{Z}(1+\theta) = x_j)|V\big].$$
(78)

The last measure we will introduce is the distribution $\nu_{1+\theta}^{\varepsilon}(\cdot|x_i^{\varepsilon})$ of $\bar{X}^{\varepsilon}(1+\theta)$ conditioned on $\bar{X}^{\varepsilon}(1)=x_i^{\varepsilon}$. We denote its atoms by $(x_j^{\varepsilon},u_{ij}^{\varepsilon})$. Thus,

$$u_{ij}^{\varepsilon} = \bar{\mathbb{P}} \big[\bar{X}^{\varepsilon} (1 + \theta) = x_j^{\varepsilon} | \bar{X}^{\varepsilon} (1) = x_i^{\varepsilon}, V \big], u_{ij} = \bar{\mathbb{P}} \big[\bar{Z} (1 + \theta) = x_j | \bar{Z} (1) = x_i, V \big].$$
 (79)

Observe that $w_{ij}^{\varepsilon} = w_i^{\varepsilon} u_{ij}^{\varepsilon}$ and $w_{ij} = w_i u_{ij}$.

Using this notation we can rewrite (75) and (76),

$$R_{\varepsilon}(\theta) = \bar{\mathbb{E}} \left[\sum_{i} w_{i}^{\varepsilon} u_{ii}^{\varepsilon} \right] \quad \text{and} \quad R_{\varepsilon}^{q}(\theta) = \bar{\mathbb{E}} \left[\sum_{i} w_{i}^{\varepsilon} \sum_{j} (u_{ij}^{\varepsilon})^{2} \right], \quad (80)$$

where the expectations are taken over all realisations of V. Obviously we have

$$R(\theta) = \bar{\mathbb{E}}\left[\sum_{i} w_i u_{ii}\right] \quad \text{and} \quad R^q(\theta) = \bar{\mathbb{E}}\left[\sum_{i,j} w_i (u_{ij})^2\right].$$
 (81)

If we prove the \mathbb{P} -a.s. convergence of the expressions inside the expectations in (80) to the corresponding expressions in (81), the proof will follow easily using the dominated convergence theorem. We want to use the results of Proposition 6.2, namely the point process convergence of ν_1^{ε} to ν_1 and $\nu_{1+\theta}^{\varepsilon}(\cdot|x_{j_i(\varepsilon)}^{\varepsilon})$ to $\nu_{1+\theta}(\cdot|x_i)$. Here, as usually, $j_i(\varepsilon)$ satisfies $(x_{j_i(\varepsilon)}, v_{j_i(\varepsilon)}) \to (x_i, v_i)$ as $\varepsilon \to 0$. Note that the point process convergence of $\nu_{1+\theta}^{\varepsilon}(\cdot|x_{j_i(\varepsilon)}^{\varepsilon})$ follows from Propositions 6.2 and 2.5(iv).

In the proof we will need one property of the atoms of different measures that is connected with Condition 1. From the point process convergence of $\bar{\mu}^{\varepsilon}$ we know that for every atom (x_l, v_l) of $\bar{\rho}$ there is a function $j_l(\varepsilon)$ such that $(x_{j_l(\varepsilon)}^{\varepsilon}, v_{j_l(\varepsilon)}^{\varepsilon})$ converges to (x_l, v_l) . From Proposition 2.5(iii) we can see that for the same function $w_{j_l(\varepsilon)}^{\varepsilon} \to w_l$, $u_{j_l(\varepsilon),j_k(\varepsilon)}^{\varepsilon} \to u_{lk}$, and thus $w_{j_l(\varepsilon),j_k(\varepsilon)}^{\varepsilon} \to w_{lk}$ as $\varepsilon \to 0$. This observation is essential, because only the point process convergence of all measures is not sufficient to imply our results.

We prove the convergence only for the quantity $R(\theta)$. The proof for $R^q(\theta)$ is entirely similar. Point process convergence, Condition 1, and the observation of the previous paragraph give

$$\sum_{i} w_{i} u_{ii} = \lim_{\varepsilon \to 0} \sum_{i} w_{j_{i}(\varepsilon)}^{\varepsilon} u_{j_{i}(\varepsilon), j_{i}(\varepsilon)}^{\varepsilon} \le \liminf_{\varepsilon \to 0} \sum_{i} w_{i}^{\varepsilon} u_{ii}^{\varepsilon}.$$
 (82)

To show the opposite bound we choose $\delta > 0$, and divide the sum in (80) into sums over three disjoint sets

$$A_{\varepsilon}(\delta) = \{i : w_i^{\varepsilon} > \delta, u_{ii}^{\varepsilon} > \delta\}$$

$$B_{\varepsilon}(\delta) = \{i : u_{ii}^{\varepsilon} \leq \delta\}$$

$$C_{\varepsilon}(\delta) = \{i : w_i^{\varepsilon} \leq \delta, u_{ii}^{\varepsilon} > \delta\}.$$
(83)

The sum over $A_{\varepsilon}(\delta)$ has necessarily finite number of terms. From point process convergence we have

$$\limsup_{\varepsilon \to 0} \sum_{i \in A_{\varepsilon}(\delta)} w_i^{\varepsilon} u_{ii}^{\varepsilon} = \sum_{i \in A(\delta)} w_i u_{ii}, \tag{84}$$

where $A(\delta)$ has the obvious meaning. For the second part we have

$$\limsup_{\varepsilon \to 0} \sum_{i \in B_{\varepsilon}(\delta)} w_i^{\varepsilon} u_{ii}^{\varepsilon} \le \delta \limsup_{\varepsilon \to 0} \sum_{i \in B_{\varepsilon}(\delta)} w_i^{\varepsilon} \le \delta, \tag{85}$$

since ν_1^{ε} is the probability measure. The last part satisfies

$$\limsup_{\varepsilon \to 0} \sum_{i \in C_{\varepsilon}(\delta)} w_i^{\varepsilon} u_{ii}^{\varepsilon} \le \limsup_{\varepsilon \to 0} \sum_{i \in C_{\varepsilon}(\delta)} w_i^{\varepsilon} \le 1 - \liminf_{\varepsilon \to 0} \sum_{i: w_i^{\varepsilon} > \delta} w_i^{\varepsilon}.$$
 (86)

The sum in the last expression has a finite number of terms. Hence

$$\limsup_{\varepsilon \to 0} \sum_{i \in C_{\varepsilon}(\delta)} w_i^{\varepsilon} u_{ii}^{\varepsilon} \le 1 - \sum_{i:w_i > \delta} w_i, \tag{87}$$

and the last sum goes to 1 as $\delta \to 0$, because ν_1 is a purely discrete measure. From (84)–(87) it is easy to see that

$$\limsup_{\varepsilon \to 0} \sum_{i} w_{i}^{\varepsilon} u_{ii}^{\varepsilon} \le \sum_{i \in A(\delta)} w_{i} u_{ii} + \delta + (1 - \sum_{i:w_{i} > \delta} w_{i})$$
 (88)

and the proof is finished by taking the limit $\delta \to 0$.

8. Proof of the sub-aging in symmetric case

We start the proof by a technical lemma that will provide the connection between the rescaled processes at time t = 1 and the process X at some large time t. Let $\varepsilon(t)$ be defined by

$$\varepsilon(t) = \sup\{\varepsilon > 0 : \varepsilon c_{\varepsilon} t \le 1\}. \tag{89}$$

We write c_t for $c_{\varepsilon(t)}$ and we define $k(t) = \varepsilon(t)c_t t$.

The next lemma defines the slowly varying function L(t) that is used in Theorem 1.3. Note that all slowly varying function that we use are slowly varying at infinity.

Lemma 8.1. (i) There exists a slowly varying function L(t) such that

$$c_t t^{\gamma} L(t) = 1. (90)$$

(ii) The function k(t) satisfies $\lim_{t\to\infty} k(t) = 1$.

The proof of this lemma is postponed to the end of the section.

The main step in proving Theorem 1.3 is the following proposition that describes the scaling of the distribution of the depth of the site where X stays at time t. We recall that

$$\gamma = \frac{\beta}{1+\beta} = \frac{1}{1+\alpha}.\tag{91}$$

Proposition 8.2. Let $F_t(u) = \mathbb{EP}(\tau(X(t))/t^{\gamma}L(t) \leq u|E)$. Then

$$\lim_{t \to \infty} F_t(u) = \mathbb{EP}(\rho(Z(1)) \le u | \rho) \equiv F(u)$$
 (92)

for all points of continuity of F(u).

We use this proposition to prove subaging for a=0.

Proof of Theorem 1.3 in the symetric case. The process X stays at the site i for an exponentially long time with mean τ_i . Using the Markov property we can write

$$\mathbb{P}[X(t') = X(t)\forall t' \in [t, t + \theta t^{\gamma} L(t)]] = \int_0^\infty e^{-\theta t^{\gamma} L(t)/u} dF_t(u/(t^{\gamma} L(t))) = \int_0^\infty e^{-\theta/u} dF_t(u).$$
(93)

By the weak convergence stated in Proposition 8.2, the last expression converges to $\int e^{-\theta/u} dF(u) = \Pi(\theta)$.

The proof of Theorem 1.3 for the asymmetric case is postponed to the next section because it is relatively complicated and relies on some notation introduced later in this section.

Proof of Proposition 8.2. We follow the similar strategy as in the proof of aging. Again we start with some notations. Let $h(\varepsilon)$ be such that $\lim_{\varepsilon\to 0} h(\varepsilon) = 1$. We write

$$\bar{\mu}^{\varepsilon}(dx) = \sum_{i \in \mathbb{Z}} c_{\varepsilon} \tau_i^{\varepsilon} \delta_{i\varepsilon}(dx) \quad \text{and} \quad \bar{\rho}(dx) = \sum_{i \in \mathbb{Z}} v_i \delta_{x_i}(dx).$$
 (94)

Similarly, the distributions of $\bar{X}^{\varepsilon}(h(\varepsilon))$ and $\bar{Z}(1)$ satisfy

$$\bar{\nu}_{h(\varepsilon)}^{\varepsilon}(dx) = \sum_{i \in \mathbb{Z}} w_i^{\varepsilon} \delta_{i\varepsilon}(dx) \quad \text{and} \quad \bar{\nu}_1(dx) = \sum_{i \in \mathbb{Z}} w_i \delta_{x_i}(dx).$$
 (95)

Here again we used the fact that the sets of positions of atoms of $\bar{\rho}$ and $\bar{\nu}_1$ are equal. Note that w_i^{ε} depends on the function h but we do not denote this dependence explicitly. We also introduce the distributions of the depth at the time $h(\varepsilon)$ resp. 1.

$$\pi_{h(\varepsilon)}^{\varepsilon}(dx) = \sum_{i \in \mathbb{Z}} w_i^{\varepsilon} \delta_{c_{\varepsilon} \tau_i^{\varepsilon}}(dx)$$
(96)

and

$$\pi_1(dx) = \sum_{i \in \mathbb{Z}} w_i \delta_{\bar{\rho}(x_i)}(dx) = \sum_{i \in \mathbb{Z}} w_i \delta_{v_i}(dx). \tag{97}$$

We claim that

Lemma 8.3.

$$\pi_{h(\varepsilon)}^{\varepsilon} \xrightarrow{v} \pi_1 \quad and \quad \pi_{h(\varepsilon)}^{\varepsilon} \xrightarrow{pp} \pi_1 \quad as \quad \varepsilon \to 0 \quad \bar{\mathbb{P}}\text{-}a.s.$$
 (98)

Proof. As usually we prove the vague convergence and Condition 1. To verify the second property, let us first observe that for any atom (v_l, w_l) of π_1 there exists x_l such that (x_l, v_l) is an atom of $\bar{\rho}$, and (x_l, w_l) is an atom of $\bar{\nu}_1$. From the point process convergences $\mu^{\varepsilon} \stackrel{pp}{\to} \bar{\rho}$, $\bar{\nu}_{h(\varepsilon)}^{\varepsilon} \stackrel{pp}{\to} \bar{\nu}_1$, and from the direct part of the Lemma 2.4 we have that for any l there exist sequences $j_l(\varepsilon)$ and $k_l(\varepsilon)$, such that $(\varepsilon j_l(\varepsilon), c_{\varepsilon} \tau_{j_l(\varepsilon)}^{\varepsilon}) \to (x_l, v_l)$ and $(\varepsilon k_l(\varepsilon), w_{k_l(\varepsilon)}^{\varepsilon}) \to (x_l, w_l)$ as $\varepsilon \to 0$. Moreover, it can be seen from Proposition 2.5(iii) that $j_l(\varepsilon) = k_l(\varepsilon)$. Putting together the last three claims and taking into account that $\bar{\mathbb{P}}$ -a.s. $x_l \neq x_m$ implies $v_l \neq v_m$, we easily show that $(c_{\varepsilon} \tau_{j_l(\varepsilon)}^{\varepsilon}, w_{j_l(\varepsilon)}^{\varepsilon}) \to (v_l, w_l)$ as $\varepsilon \to 0$.

We should now verify the vague convergence. Let f be a nonnegative, continuous function with compact support. We use I_{δ} to denote the open rectangle $(-\delta^{-1}, \delta^{-1}) \times (\delta, 2)$. By (96) we have

$$\int f(x)\pi_{h(\varepsilon)}^{\varepsilon}(dx) = \sum_{i\in\mathbb{Z}} w_{i}^{\varepsilon} f(c_{\varepsilon}\tau_{i\varepsilon}^{\varepsilon}) =$$

$$= \sum_{i:(i\varepsilon,w_{i}^{\varepsilon})\in I_{\delta}} w_{i}^{\varepsilon} f(c_{\varepsilon}\tau_{i\varepsilon}^{\varepsilon}) + \sum_{i:(i\varepsilon,w_{i}^{\varepsilon})\notin I_{\delta}} w_{i}^{\varepsilon} f(c_{\varepsilon}\tau_{i\varepsilon}^{\varepsilon}).$$
(99)

From the point process convergence of $\bar{\nu}_{h(\varepsilon)}^{\varepsilon}$ we know that for all but countably many $\delta > 0$ and for ε large enough the number of atoms of $\bar{\nu}_{h(\varepsilon)}^{\varepsilon}$ in I_{δ} is finite and is equal to the number of atoms of $\bar{\nu}_1$ in I_{δ} . Moreover, by the first part of Lemma 2.4 we have for any such atom (x_l, w_l) the sequence of atoms $(\varepsilon j_l(\varepsilon), w_{j_l(\varepsilon)}^{\varepsilon})$ converging to (x_l, w_l) . By the same reasoning as in the previous paragraph the sequence $c_{\varepsilon}\tau_{j_l(\varepsilon)}^{\varepsilon}$ converges as $\varepsilon \to 0$ to $\bar{\rho}(x_l) = v_l$. Thus, by continuity of f we have

$$\lim_{\varepsilon \to 0} \sum_{i:(i\varepsilon, w_i^{\varepsilon}) \in I_{\delta}} w_i^{\varepsilon} f(c_{\varepsilon} \tau_{i\varepsilon}^{\varepsilon}) = \sum_{i:(x_i, w_i) \in I_{\delta}} w_i f(v_i).$$
 (100)

The right hand side of the last equation is bounded by $||f||_{\infty}$ and increases as δ decreases. Thus, its limit as $\delta \to 0$ exist and is equal to $\int f(x)\pi_1(dx)$.

The second sum in (99) is bounded by

$$C\sum_{i:(i\varepsilon,w_i^{\varepsilon})\notin I_{\delta}} w_i^{\varepsilon} = C\left(1 - \sum_{i:(i\varepsilon,w_i^{\varepsilon})\in I_{\delta}} w_i^{\varepsilon}\right). \tag{101}$$

Using the same argument as in (87) we have

$$\lim_{\delta \to 0} \limsup_{\varepsilon \to 0} \left(1 - \sum_{i: (i\varepsilon, w_i^{\varepsilon}) \in I_{\delta}} w_i^{\varepsilon} \right) = \lim_{\delta \to 0} \left(1 - \sum_{i: (x_i, w_i) \in I_{\delta}} w_i \right) = 0, \quad (102)$$

since the finite time distribution of \bar{Z} is discrete.

We can now finish the proof of Proposition 8.2. By definition (33) of $X^{\varepsilon}(t)$ we have

$$F_t(u) = \mathbb{P}[\tau(X(t))/t^{\gamma}L(t) \le u] = \mathbb{P}[\tau(\varepsilon^{-1}X^{\varepsilon}(t\varepsilon c_{\varepsilon}))/t^{\gamma}L(t) \le u].$$
(103)

Inserting the definition (34) of μ^{ε} into the last claim yields

$$F_t(u) = \mathbb{P}[c_{\varepsilon}^{-1} \mu^{\varepsilon}(X^{\varepsilon}(t \varepsilon c_{\varepsilon}))/t^{\gamma} L(t) \le u]. \tag{104}$$

Setting $\varepsilon = \varepsilon(t)$ and using the equality of the distributions $\bar{X}^{\varepsilon} \sim X^{\varepsilon}$, $\bar{\mu}^{\varepsilon} \sim \mu^{\varepsilon}$, and Lemma 8.1, we get

$$F_t(u) = \bar{\mathbb{P}}[\bar{\mu}^{\varepsilon(t)}(\bar{X}^{\varepsilon(t)}(k(t))) \le u]. \tag{105}$$

By definition (96) of $\pi_{k(t)}^{\varepsilon(t)}$ we have

$$1 - F_t(u) = \bar{\mathbb{E}}\bar{\mathbb{P}}[\bar{\mu}^{\varepsilon(t)}(\bar{X}^{\varepsilon(t)}(k(t))) > u|V] = \bar{\mathbb{E}}\Big[\sum_{i: c_t \tau_i^{\varepsilon(t)} > u} w_i^{\varepsilon(t)}\Big]. \quad (106)$$

The point process convergence proved in Lemma 8.3 implies that the sum in the last expectation converges $\bar{\mathbb{P}}$ -a.s. for all u such that $u \neq v_i$ for all i.

$$\lim_{t \to \infty} \sum_{i: c_t \tau_i^{\varepsilon(t)} > u} w_i^{\varepsilon(t)} = \sum_{i: v_i > u} w_i = \bar{\mathbb{P}}[\bar{\rho}(\bar{Z}(1)) > u | V]$$
 (107)

Using the fact that (ρ, Z) has the same distribution as $(\bar{\rho}, Z)$ and applying dominated convergence theorem it is easy to finish the proof.

Proof of Lemma 8.1. Let $L_1(t)$ be defined by

$$\mathbb{P}[\tau_0 > t] = t^{-\alpha} L_1(t). \tag{108}$$

Since τ_0 is in the domain of attraction of the stable law with index α , the function L_1 is slowly varying.

We first show the second claim of the lemma, namely that $k(t) \to 1$ as $t \to \infty$. It is easy to see from (89) that $k(t) \ge 1$. To get an upper bound take $\delta > 0$ and assume that

$$\lim_{t \to \infty} \sup k(t) = \lim_{t \to \infty} \sup \varepsilon(t) c_t t \ge 1 + \delta. \tag{109}$$

If this is true, then there is a sequence t_n such that $t_n \to \infty$ as $n \to \infty$, and $\varepsilon(t_n)c_{t_n}t_n \geq (1+\delta)$. Using again (89) we get

$$\varepsilon(t_n)c_{t_n} \ge (1+\delta)t_n^{-1} \ge (1+\delta)\lim_{\varepsilon\uparrow\varepsilon(t_n)}\varepsilon c_{\varepsilon}$$
 (110)

This means that $(1+\delta)c_{\varepsilon(t_n)}^{-1} \leq \lim_{\varepsilon \uparrow \varepsilon(t_n)} c_{\varepsilon}^{-1}$. Using the definition (49) of c_{ε} , it is easy to see that this can only happen if there is a sequence s_n such that $s_n \to \infty$ as $n \to \infty$, and $\mathbb{P}[\tau_0 > s_n] = \mathbb{P}[\tau_0 > (1+\delta)s_n]$. However, then

$$\frac{L_1((1+\delta)s_n)}{L_1(s_n)} = \frac{(1+\delta)^{\alpha} s_n^{\alpha} \mathbb{P}[\tau_0 > (1+\delta)s_n]}{s_n^{\alpha} \mathbb{P}[\tau_0 > s_n]} = (1+\delta)^{\alpha}$$
(111)

and this leads to contradiction since L_1 is a slowly varying function. Therefore (109) is false and the second part of the lemma is proved.

To verify the first claim of the lemma we should only prove that L(t) is slowly varying. From definition (49) of c_{ε} we get

$$\varepsilon^{-1}\mathbb{P}[\tau_0 > c_{\varepsilon}^{-1}] \to 1 \quad \text{as} \quad \varepsilon \to 0.$$
 (112)

Indeed, it is easy to see that $\varepsilon^{-1}\mathbb{P}[\tau_0 > c_{\varepsilon}^{-1}] \leq 1$. Take $\eta > 0$, the lower bound follows from

$$(1+2\eta)^{-\alpha} = \lim_{\varepsilon \to 0} \frac{\mathbb{P}[\tau_0 > \frac{1+2\eta}{1+\eta} c_{\varepsilon}^{-1}]}{\mathbb{P}[\tau_0 > \frac{1}{1+\eta} c_{\varepsilon}^{-1}]} \le \liminf_{\varepsilon \to 0} \varepsilon^{-1} \mathbb{P}[\tau_0 > c_{\varepsilon}^{-1}]$$
(113)

since η is arbitrary. From (112) and (108) we get

$$\varepsilon^{-1}c_{\varepsilon}^{\alpha}L_1(c_{\varepsilon}^{-1}) \to 1 \quad \text{as} \quad \varepsilon \to 0.$$
 (114)

Using (114) and $k(t) \rightarrow 1$ we get

$$c_t t^{\gamma} L_1^{\gamma}(c_t^{-1}) \to 1 \quad \text{as} \quad t \to \infty.$$
 (115)

We want to show that $c_t = t^{-\gamma}L(t)^{-1}$ where L(t) is slowly varying. Choose k > 0 and define $d_t = L(t)/L(kt)$. Take $\eta > 0$ small and assume that $\lim \inf_{t\to\infty} d_t < 1 - 2\eta$. We choose $\delta > 0$ and we consider

t large enough such that $c_t t^{\gamma} L_1^{\gamma}(c_t^{-1}) \in (1 - \delta, 1 + \delta)$. This can be done by (115). We have

$$d_t = \frac{L(t)}{L(kt)} = \frac{c_{kt}}{c_t} k^{\gamma} \ge \frac{1 - \delta}{1 + \delta} \cdot \frac{L_1^{\gamma}(c_t^{-1})}{L_1^{\gamma}(c_{kt}^{-1})} = \frac{1 - \delta}{1 + \delta} \cdot \frac{L_1^{\gamma}(c_t^{-1})}{L_1^{\gamma}(d_t^{-1}c_t^{-1}k^{\gamma})}.$$
(116)

Our assumption implies that there exists a sequence t_n such that $d_{t_n}^{-1} > 1 + \eta$ for all n. Since L_1 is slowly varying, we know that for arbitrary $\theta > 0$ there exists x_0 such that for all $l > 1 + \eta$ and $x > x_0$ we have $L_1(lx) \leq l^{\theta}L_1(x)$. This implies that for n large enough we have

$$d_{t_n} \ge \frac{1 - \delta}{1 + \delta} \cdot \frac{L_1^{\gamma}(c_{t_n}^{-1})}{d_{t_n}^{-\gamma\theta} L_1^{\gamma}(c_{t_n}^{-1}k^{\gamma})}.$$
 (117)

Taking the limit $n \to \infty$, using that $c_{t_n} \to \infty$ and that L_1 is slowly varying we get

$$\liminf_{n \to \infty} d_{t_n}^{1+\gamma\theta} \ge \frac{1-\delta}{1+\delta}.$$
(118)

For every η we can take δ and θ such that the last equation is in contradiction with $\liminf_{t\to\infty} d_t < 1 - 2\eta$. Thus $\liminf_{t\to\infty} d_t \geq 1$. The proof of the upper bound follows from

$$d_{t_n} \le \frac{1+\delta}{1-\delta} \cdot \frac{d_{t_n}^{\gamma \theta} L_1^{\gamma} (k^{-\gamma} c_{kt_n}^{-1})}{L_1^{\gamma} (c_{kt_n}^{-1})}.$$
 (119)

This can be proved if one assumes that $\limsup_{t\to\infty} d_t \ge 1 + 2\eta$ and it leads to a contradiction similarly as in (118).

9. Proof of sub-aging in the non-symmetric case

If a > 0, the jump rates depend also on the depths of the neighbouring sites. As it is easy to see from the definition of τ_i^{ε} , the depth of the neighbouring sites of some very deep trap does not converge $\bar{\mathbb{P}}$ -a.s. (By very deep trap we mean here a trap where X has a large chance to stay at time t.) On the other hand, we expect (see Rinn, Maass, and Bouchaud, 2000) that the depth of these sites is, at least if t_w is large, almost independent of the diffusion and has the same distribution as E_0 . We will show that this expectation is correct.

We consider the function $\Pi(t, t + f_a(t, \theta))$. By its definition we have

$$\Pi(t, t + f_a(t, \theta)) = \mathbb{E}\left[\sum_{i \in \mathbb{Z}} \mathbb{P}(X(t) = i|E) \exp(-(w_{i,i+1} + w_{i,i-1})f_a(t, \theta))\right].$$
(120)

The rates $w_{i,i+1}$ and $w_{i,i-1}$ can be expressed using the variables τ_i

$$w_{i,i+1} + w_{i,i-1} = \frac{\tau_{i-1}^a + \tau_{i+1}^a}{\tau_i^{1-a}} \left\lceil \frac{\mathbb{E}(\exp(-2a\beta E_0))}{2} \right\rceil^{1-2a}.$$
 (121)

We use K to denote the constant in the brackets in the last expression. Then, taking $\varepsilon = \varepsilon(t)$ as in (89),

$$\Pi(t, t + f_a(t, \theta)) = \mathbb{E}\left[\sum_{i \in \mathbb{Z}} \mathbb{P}(X^{\varepsilon}(t\varepsilon c_{\varepsilon}) = \varepsilon i | E) \exp\left(-Kf_a(t, \theta) \frac{\tau_{i+1}^a + \tau_{i-1}^a}{\tau_i^{1-a}}\right)\right] \\
= \overline{\mathbb{E}}\left[\sum_{i \in \mathbb{Z}} w_i(t) \exp\left(-Kf_a(t, \theta) \frac{(\tau_{i+1}^{\varepsilon(t)})^a + (\tau_{i-1}^{\varepsilon(t)})^a}{(\tau_i^{\varepsilon(t)})^{1-a}}\right)\right], \tag{122}$$

where $w_i(t) = \bar{\mathbb{P}}(\bar{X}^{\varepsilon(t)}(t\varepsilon(t)c_t) = i\varepsilon(t)|V) = \bar{\mathbb{P}}(\bar{X}^{\varepsilon(t)}(k(t)) = i\varepsilon(t)|V)$. Let m > 0 large and $\eta > 0$ small. We use $J_m^{\eta} = J_m^{\eta}(V)$ to denote the set of deep traps not far from the origin

$$J_m^{\eta} = \{ x \in [-m, m] : V(x) - V(x-) \ge \eta \}. \tag{123}$$

Let $T_m^{\eta}(\varepsilon)$ be the set of sites corresponding to J_m^{η} at the scale ε

$$T_m^{\eta}(\varepsilon) = \{ i \in \mathbb{Z} : (i\varepsilon, (i+1)\varepsilon] \cap J_m^{\eta} \neq \emptyset \}. \tag{124}$$

Note that J_m^{η} and $T_m^{\eta}(\varepsilon)$ are $\bar{\mathbb{P}}$ -a.s. finite sets.

In the following proposition we show that it is possible to choose m and η such that $\bar{X}^{\varepsilon(t)}(k(t))$ is with an arbitrarily large probability in $T_m^{\eta}(\varepsilon(t))$. This can be regarded as a stronger version of the localisation effect (8) since the size of the set $T_m^{\eta}(\varepsilon)$ can be bounded uniformly in ε by $|J_m^{\eta}|$.

Proposition 9.1. Let $h(\varepsilon)$ be such that $\lim_{\varepsilon \to 0} h(\varepsilon) = 1$. Then for every $\delta > 0$ there exist m, η , and ε_0 such that for $\varepsilon < \varepsilon_0$

$$\bar{\mathbb{P}}\big[\bar{\mathbb{P}}(\varepsilon^{-1}\bar{X}^{\varepsilon}(h(\varepsilon)) \in T_m^{\eta}(\varepsilon)|V) > 1 - \delta\big] > 1 - \delta. \tag{125}$$

We postpone the proof of this proposition to the end of this section and we use it to further simplify (122). Let $\delta > 0$ and let m and η be such that (125) holds. We divide the sum in (122) into two parts. The contribution of the sum over $i \notin T_m^{\eta}(\varepsilon)$ is not important. Indeed, by Proposition 9.1, for all t large enough

$$\bar{\mathbb{E}}\left[\sum_{i\in\mathbb{Z}\backslash T_m^{\eta}(\varepsilon(t))} w_i(t) \exp\left(-f_a(t,\theta)K \frac{(\tau_{i+1}^{\varepsilon(t)})^a + (\tau_{i-1}^{\varepsilon(t)})^a}{(\tau_i^{\varepsilon(t)})^{1-a}}\right)\right] \\
\leq \bar{\mathbb{E}}\left[\sum_{i\in\mathbb{Z}\backslash T_m^{\eta}(\varepsilon(t))} w_i(t)\right] \leq 2\delta. \quad (126)$$

To estimate the contribution of the sum over $i \in T_m^{\eta}(\varepsilon)$ we define the set of neighbours of deep sites

$$N_m^{\eta}(\varepsilon) = \{ i \in \mathbb{Z} \setminus T_m^{\eta}(\varepsilon) : \exists j \in T_m^{\eta}(\varepsilon) \text{ such that } |i - j] = 1 \}.$$
 (127)

Let $\dot{\sigma}_i^{\varepsilon}$ be a sequence of i.i.d. random variables defined on $\bar{\Omega}$ that are independent of V and have the same distribution as τ_0^{ε} conditioned on $J_m^{\eta} \cap (0, \varepsilon] = \emptyset$. Let $\hat{\sigma}_i^{\varepsilon} = \min(\dot{\sigma}_i^{\varepsilon}, c_{\varepsilon}^{-1/2})$. We define

$$\dot{\tau}_{i}^{\varepsilon} = \begin{cases} \dot{\sigma}_{i}^{\varepsilon} & \text{for } i \in N_{m}^{\eta}(\varepsilon) \\ \tau_{i}^{\varepsilon} & \text{otherwise.} \end{cases} \quad \text{and} \quad \hat{\tau}_{i}^{\varepsilon} = \begin{cases} \hat{\sigma}_{i}^{\varepsilon} & \text{for } i \in N_{m}^{\eta}(\varepsilon) \\ \tau_{i}^{\varepsilon} & \text{otherwise.} \end{cases}$$
(128)

We define measures $\dot{\mu}^{\varepsilon}$, $\hat{\mu}^{\varepsilon}$ and scaling functions \dot{S}_{ε} , \hat{S}_{ε} similarly as in (55) and (57) but using $\dot{\tau}_{i}^{\varepsilon}$, $\hat{\tau}_{i}^{\varepsilon}$ instead of τ_{i}^{ε} . Further, let

$$\dot{X}^{\varepsilon}(t) = X(\dot{\mu}^{\varepsilon}, \varepsilon \dot{S}_{\varepsilon}(\varepsilon^{-1} \cdot))(t) \text{ and } \hat{X}^{\varepsilon}(t) = X(\hat{\mu}^{\varepsilon}, \varepsilon \hat{S}_{\varepsilon}(\varepsilon^{-1} \cdot))(t), \quad (129)$$

and let $\dot{w}_i(t)$, $\hat{w}_i(t)$ be defined similarly as $w_i(t)$.

To finish the proof of the theorem we will need four technical lemmas.

Lemma 9.2. For every fixed realisation of $\dot{\sigma}_i^{\varepsilon}$, $\bar{\mathbb{P}}$ -a.s.

$$\varepsilon \hat{S}_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \hat{\mu}^{\varepsilon} \xrightarrow{v} \bar{\rho} \quad and \quad \varepsilon \hat{S}_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \hat{\mu}^{\varepsilon} \xrightarrow{pp} \bar{\rho} \quad as \ \varepsilon \to 0.$$
 (130)

Therefore, the distribution of $\hat{X}^{\varepsilon(t)}(k(t))$ converges as $t \to \infty$ weakly and in the point process sense to the distribution of $\bar{Z}(1)$.

In particular, for all $x \in J_m^{\eta}$

$$\lim_{t \to \infty} \hat{w}_{j_{\varepsilon(t)}}(t) = \bar{\mathbb{P}}(\bar{Z}(1) = x | V) \equiv w_x \quad and \quad \lim_{\varepsilon \to 0} c_{\varepsilon} \hat{\tau}_{j_{\varepsilon}}^{\varepsilon} = \bar{\rho}(x).$$

$$(131)$$
where $j_{\varepsilon} = j_{\varepsilon}(x) \in T_m^{\eta}(\varepsilon)$ satisfies $x \in (\varepsilon j_{\varepsilon}, \varepsilon (j_{\varepsilon} + 1)].$

Proof. The proof of the first part of this lemma is very similar to the proofs of Lemma 5.2 and Proposition 5.1, the finite number of changes of neighbours of the deep traps looses its influence as $\varepsilon \to 0$. Therefore, we only describe modifications that must be done in the original proofs.

To get an equivalent of Lemma 5.2 we must show that $\varepsilon \hat{S}_{\varepsilon}(\lfloor \varepsilon^{-1}y \rfloor) = \varepsilon \sum_{j=0}^{\lfloor \varepsilon^{-1}y \rfloor} \hat{r}_{j}^{\varepsilon}$ converges to y. Since J_{m}^{η} is finite, only the finite number of $\hat{r}_{j}^{\varepsilon}$'s is influenced by changing the sequence of τ 's. Since $\hat{r}_{j}^{\varepsilon}$ are bounded, the contribution of the changed part of the sum tends to zero as $\varepsilon \to 0$. The rest of the sum can be treated in the same way as in the proof of Lemma 5.2.

Further, we must show the vague convergence and Condition 1 for the measures $\varepsilon \hat{S}_{\varepsilon}(\varepsilon^{-1}\cdot) \circ \hat{\mu}^{\varepsilon}$. Let x be position of an atom of ρ and let $i(\varepsilon)$ be given by $x \in (\varepsilon i(\varepsilon), \varepsilon (i(\varepsilon) + 1)]$. It is easy to observe that $i(\varepsilon) \notin N_m^{\eta}(\varepsilon)$ for all ε small enough. The proof of Condition 1 can be then finished using the same reasoning as before.

To show the vague convergence let f(x) be a bounded continuous function with bounded support. Then

$$\left| \int f(x) (\varepsilon \hat{S}_{\varepsilon}(\varepsilon^{-1} \cdot) \circ \hat{\mu}^{\varepsilon}) (dx) - \int f(x) (\varepsilon S_{\varepsilon}(\varepsilon^{-1} \cdot) \circ \bar{\mu}^{\varepsilon}) (dx) \right|$$

$$\leq \left| \sum_{i \in N_{m}^{\eta}(\varepsilon)} f(\varepsilon S_{\varepsilon}(i)) g_{\varepsilon} (\bar{\rho}(i\varepsilon, (i+1)\varepsilon)) \right| + \left| \sum_{i \in N_{m}^{\eta}(\varepsilon)} f(\varepsilon \hat{S}_{\varepsilon}(i)) c_{\varepsilon} \hat{\sigma}_{i}^{\varepsilon} \right|.$$

$$(132)$$

The contribution of the first term can be proved to be small observing that J^{ε}_{δ} (defined in (63)) satisfies $J^{\varepsilon}_{\delta} \cap N^{\eta}_{m}(\varepsilon) = \emptyset$ for all $\delta > 0$ if ε is small enough. The second term in (132) is also negligible since $\hat{\sigma}^{\varepsilon}_{i} \leq c^{-1/2}_{\varepsilon}$ and $|N^{\eta}_{m}(\varepsilon)| \leq 2|J^{\eta}_{m}|$ is a.s. finite.

The convergence of $\hat{X}^{\varepsilon(t)}(k(t))$ and of $w_{j_{\varepsilon(t)}}(t)$ is then a consequence of the first part of the lemma and Proposition 2.5.

Lemma 9.3. The sequence $\dot{\tau}_i^{\varepsilon}$ has the same distribution as τ_i .

Proof. The proof is obvious because the distribution of $\dot{\sigma}_i^{\varepsilon}$ is chosen to be equal to the distribution of τ_i^{ε} conditioned on $i \notin T_m^{\eta}(\varepsilon)$.

Lemma 9.4. For $\overline{\mathbb{P}}$ -a.e. realisation of V

$$\lim_{\varepsilon \to 0} \bar{\mathbb{P}}(\exists i \in N_m^{\eta}(\varepsilon) : \dot{\sigma}_i^{\varepsilon} \neq \hat{\sigma}_i^{\varepsilon} | V) = 0.$$
 (133)

Proof. The probability that $\dot{\sigma}_i^{\varepsilon} \geq c_{\varepsilon}^{-1/2}$ tends to zero. Since $N_m^{\eta}(\varepsilon)$ is a.s. finite, the proof is finished.

Lemma 9.5. As $\varepsilon \to 0$ the random variables $\hat{\sigma}_0^{\varepsilon}$ converge weakly to τ_0 .

Proof. By definition of $\dot{\sigma}_i$, $\bar{\mathbb{P}}(\dot{\sigma}_0^{\varepsilon} \leq a) = \bar{\mathbb{P}}(\tau_0^{\varepsilon} \leq a | J_m^{\eta} \cap (0, \varepsilon] = \emptyset)$. Since the probability of the conditioning event tends to one as $\varepsilon \to 0$ and τ_0^{ε} has the same distribution as τ_0 , this converges to $\bar{\mathbb{P}}(\tau_0 \leq a)$. Therefore, $\dot{\sigma}_0^{\varepsilon}$ converges weakly to τ_0 . Since $\hat{\sigma}_0^{\varepsilon} = \min(\dot{\sigma}_0^{\varepsilon}, c_{\varepsilon}^{-1/2})$ and $c_{\varepsilon}^{-1/2} \to \infty$ as $\varepsilon \to 0$, the lemma follows.

We can now estimate the contribution of the sum over $i \in T_m^{\eta}(\varepsilon)$ in (122). Using Lemma 9.3 we get

$$\bar{\mathbb{E}}\left[\sum_{i \in T_m^{\eta}(\varepsilon(t))} w_i(t) \exp\left(-Kf_a(t,\theta) \frac{(\tau_{i+1}^{\varepsilon(t)})^a + (\tau_{i-1}^{\varepsilon(t)})^a}{(\tau_i^{\varepsilon(t)})^{1-a}}\right)\right]$$

$$= \bar{\mathbb{E}}\left[\sum_{i \in T_m^{\eta}(\varepsilon(t))} \dot{w}_i(t) \exp\left(-Kf_a(t,\theta) \frac{(\dot{\tau}_{i+1}^{\varepsilon(t)})^a + (\dot{\tau}_{i-1}^{\varepsilon(t)})^a}{(\dot{\tau}_i^{\varepsilon(t)})^{1-a}}\right)\right]$$

$$= \bar{\mathbb{E}}\left[\sum_{i \in T_m^{\eta}(\varepsilon(t))} \dot{w}_i(t) \exp\left(-Kf_a(t,\theta) \frac{(\hat{\tau}_{i+1}^{\varepsilon(t)})^a + (\hat{\tau}_{i-1}^{\varepsilon(t)})^a}{(\hat{\tau}_i^{\varepsilon(t)})^{1-a}}\right)\right] + R_1(\varepsilon(t))$$
(134)

The error term $R_1(\varepsilon)$ can be bounded by

$$|R_1(\varepsilon)| \le \bar{\mathbb{E}}[\bar{\mathbb{P}}(\exists i \in N_m^{\eta}(\varepsilon) : \dot{\sigma}_i^{\varepsilon} \ne \hat{\sigma}_i^{\varepsilon}|V)] \to 0 \quad \text{as } \varepsilon \to 0$$
 (135)

by Lemma 9.4 and the dominated convergence theorem. Recall that $f_a(t,\theta) = \theta t^{\gamma(1-\alpha)} L(t)^{1-\alpha}$. Therefore, using Lemma 8.1, the main term in (134) can be rewritten as

$$\bar{\mathbb{E}}\Big[\sum_{x\in J_m^n} w_x \exp\Big(-Kf_a(t,\theta) \frac{(\hat{\tau}_{j_{\varepsilon(t)}(x)+1}^{\varepsilon(t)})^a + (\hat{\tau}_{j_{\varepsilon(t)}(x)-1}^{\varepsilon(t)})^a}{c_t^{a-1}\bar{\rho}(x)^{1-a}}\Big)\Big] + R_2(\varepsilon(t))$$

$$= \bar{\mathbb{E}} \left[\sum_{x \in J_m^{\eta}} w_x \exp\left(-K\theta \frac{(\hat{\sigma}_{j_{\varepsilon(t)}(x)+1}^{\varepsilon(t)})^a + (\hat{\sigma}_{j_{\varepsilon(t)}(x)-1}^{\varepsilon(t)})^a}{\bar{\rho}(x)^{1-a}} \right) \right] + R_2(\varepsilon(t)), \tag{136}$$

where $j_{\varepsilon}(x)$ is defined as in Lemma 9.2 and $R_2(\varepsilon)$ is an error that comes from the replacement of $\hat{w}_i(t)$ and $\hat{\tau}_i^{\varepsilon(t)}$ by w_x and $c_t\bar{\rho}(x)$. It follows from Lemma 9.2 that $|R_2(\varepsilon)| \to 0$ as $\varepsilon \to 0$.

We can now easily compute the expectation over $\hat{\sigma}$ in (136). Let $g_a^{\varepsilon}(\lambda)$ denote the Laplace transform of $K(\hat{\sigma}_0^{\varepsilon})^a$, $g_a^{\varepsilon}(\lambda) = \bar{\mathbb{E}}(\exp(-\lambda K(\hat{\sigma}_0^{\varepsilon})^a))$, and let $g_a(\lambda) = \mathbb{E}(\exp(-\lambda K\tau_0^a))$. Since τ_0 has the same distribution as $\exp(\beta E_0)\mathbb{E}(\exp(-2a\beta E_0))/2$, $K\tau_0^a$ has the same distribution as

$$2^{a-1} \exp(a\beta E_0) \left(\mathbb{E}(\exp(-2a\beta E_0)) \right)^{1-a} \equiv T_a, \tag{137}$$

and $g_a(\lambda) = \mathbb{E}(e^{-\lambda T_a})$ as required by Theorem 1.3. From Lemma 9.5 it follows that $\lim_{\varepsilon \to 0} g_a^{\varepsilon}(\lambda) = g_a(\lambda)$. Using this notation, (126), (134), and (136) we get

$$\lim_{t \to \infty} \sup \Pi(t, f_a(t, \theta))$$

$$\leq \lim_{\varepsilon \to 0} \sup \bar{\mathbb{E}} \left[\sum_{x \in J_m^{\eta}} w_x \exp\left(-K\theta \frac{(\hat{\sigma}_{j_{\varepsilon}(x)+1}^{\varepsilon})^a + (\hat{\sigma}_{j_{\varepsilon}(x)-1}^{\varepsilon})^a}{\bar{\rho}(x)^{1-a}}\right) \right] + 2\delta$$

$$= \bar{\mathbb{E}} \left[\sum_{x \in J_m^{\eta}} w_x g_a^2 (\theta \bar{\rho}(x)^{a-1}) \right] + 2\delta.$$
(138)

Inserting the remaining atoms of $\bar{\rho}$ inside the sum, making again an error of order at most 2δ , we get

$$\limsup_{t \to \infty} \Pi(t, f_a(t, \theta)) \le \bar{\mathbb{E}} \left[\sum_x w_x g_a^2(\theta \bar{\rho}(x)^{a-1}) \right] + 4\delta.$$
 (139)

An analogous calculation gives

$$\liminf_{t \to \infty} \Pi(t, f_a(t, \theta)) \ge \bar{\mathbb{E}} \left[\sum_{x} w_x g_a^2(\theta \bar{\rho}(x)^{a-1}) \right] - 4\delta.$$
 (140)

Since δ was arbitrary we have

$$\Pi(\theta) = \int_0^\infty g_a^2(\theta u^{a-1}) dF(u), \tag{141}$$

which finishes the proof of sub-aging in the asymmetric situation. We still have to show Proposition 9.1

Proof of Proposition 9.1. The claim follows from the existence of η and m such that

$$\bar{\mathbb{P}}[\bar{\mathbb{P}}(\bar{Z}(1) \in J_m^{\eta}|V) \ge 1 - \delta/2] \ge 1 - \delta/2, \tag{142}$$

and from the \mathbb{P} -a.s. point process convergence of the distribution of $\hat{X}^{\varepsilon}(1)$ to that of $\bar{Z}(1)$. Namely, for $\bar{\mathbb{P}}$ -a.e. realisation of V it follows from Proposition 2.5(iii,iv) that there is $\varepsilon(V)>0$ such that for $\varepsilon<\varepsilon(V)$

$$\left|\bar{\mathbb{P}}(\bar{Z}(1) \in J_m^{\eta}|V) - \bar{\mathbb{P}}(\bar{X}^{\varepsilon}(h(\varepsilon)) \in T_m^{\eta}(\varepsilon)|V)\right| \le \delta/2. \tag{143}$$

We then take ε_0 such that $\bar{\mathbb{P}}(\varepsilon(V) > \varepsilon_0) > 1 - \delta/2$.

We should still verify (142). It is equivalent to

$$\bar{\mathbb{P}}[\bar{\mathbb{P}}(\bar{Z}(1) \notin J_m^{\eta}|V) \le \delta/2] \ge 1 - \delta/2. \tag{144}$$

The last claim can be easily verified if we show

$$\bar{\mathbb{P}}[\bar{Z}(1) \notin J_m^{\eta}] = \bar{\mathbb{E}}[\bar{\mathbb{P}}(\bar{Z}(1) \notin J_m^{\eta}|V)] \le \delta^2/4. \tag{145}$$

Indeed, assume that (144) is not true, i.e.

$$\bar{\mathbb{P}}[\bar{\mathbb{P}}(\bar{Z}(1) \notin J_m^{\eta}|V) > \delta/2] > \delta/2. \tag{146}$$

Then clearly

$$\bar{\mathbb{E}}[\bar{\mathbb{P}}(\bar{Z}(1) \notin J_m^{\eta}|V)] > \delta^2/4, \tag{147}$$

in contradiction with (145).

We establish claim (145) using two lemmas.

Lemma 9.6. Let $\eta(t) = t^{1/(1+\alpha)}$ and $m(t) = t^{\alpha/(1+\alpha)}$. Then

$$\bar{\mathbb{P}}(\bar{Z}(1) \in J_m^{\eta}) = \bar{\mathbb{P}}(\bar{Z}(t) \in J_{m(t)}^{\eta(t)}). \tag{148}$$

Lemma 9.7. For every δ' there exist m' and η' such that

$$\int_{0}^{1} \bar{\mathbb{P}}(\bar{Z}(t) \in J_{m'}^{\eta'}) dt \ge 1 - \delta'.$$
 (149)

We first finish the proof of Proposition 9.1. The Lemma 9.7 ensures the existence of $t \in (0,1)$ such that $\bar{\mathbb{P}}(Z(t) \in J_{m'}^{\eta'}) \geq 1 - \delta'$. The claim (145) then follows from Lemma 9.6, choosing $\delta' = \delta^2/4$, $m = t^{-\alpha/(1+\alpha)}m'$, and $\eta = t^{-1/(1+\alpha)}\eta'$.

Proof of Lemma 9.6. The pair

$$(W_{\lambda}(t), V_{\lambda}(x)) \equiv \left(\lambda W(\lambda^{-2}t), \lambda^{1/\alpha} V(\lambda^{-1}x)\right) \tag{150}$$

has the same distribution as (W(t), V(x)). The measure $\bar{\rho}_{\lambda}$ associated to V_{λ} can be written as

$$\bar{\rho}_{\lambda} = \sum_{x_i} (V_{\lambda}(x_i) - V_{\lambda}(x_i)) \delta_{x_i} = \lambda^{1/\alpha} \sum_{y_i} (V(y_i) - V(y_i)) \delta_{\lambda y_i}. \tag{151}$$

We thus have

$$\phi_{\lambda}(t) \equiv \int \ell_{\lambda}(t, y) \bar{\rho}_{\lambda}(dy) = \int \lambda \ell(\lambda^{-2}t, \lambda^{-1}y) \bar{\rho}_{\lambda}(dy)$$

$$= \sum_{y_i} \lambda \ell(\lambda^{-2}t, y_i) \lambda^{1/\alpha} (V(y_i) - V(y_i - y_i)) = \lambda^{(\alpha + 1)/\alpha} \phi(\lambda^{-2}t)$$
(152)

and therefore its generalised inverse satisfies $\psi_{\lambda}(t) = \lambda^2 \psi(\lambda^{-(\alpha+1)/\alpha}t)$. The rescaled singular diffusion defined by $\bar{Z}_{\lambda} = W_{\lambda}(\psi_{\lambda}(t))$ that has the same distribution as \bar{Z} thus satisfies

$$\bar{Z}_{\lambda}(t) = W_{\lambda}(\psi_{\lambda}(t)) = \lambda \bar{Z}(\lambda^{-(\alpha+1)/\alpha}t).$$
 (153)

Clearly, the triplet $(W_{\lambda}, V_{\lambda}, \bar{Z}_{\lambda})$ has the same distribution as (W, V, \bar{Z}) too. We thus have

$$\bar{\mathbb{P}}(\bar{Z}(1) \in J_m^{\eta}(V)) = \bar{\mathbb{P}}(\bar{Z}_{\lambda}(1) \in J_m^{\eta}(V_{\lambda})). \tag{154}$$

The set $J_m^{\eta}(V_{\lambda})$ satisfies $J_m^{\eta}(V_{\lambda}) = \lambda J_{m\lambda^{-1}}^{\eta\lambda^{-1/\alpha}}(V)$ as can be easily verified from the scaling of V or from (151) and thus

$$\bar{\mathbb{P}}(\bar{Z}(1) \in J_m^{\eta}(V)) = \bar{\mathbb{P}}(\lambda \bar{Z}(\lambda^{-(\alpha+1)/\alpha}) \in \lambda J_{m\lambda^{-1}}^{\eta\lambda^{-1/\alpha}}(V)) =
= \bar{\mathbb{P}}(\bar{Z}(\lambda^{-(\alpha+1)/\alpha}) \in J_{m\lambda^{-1}}^{\eta\lambda^{-1/\alpha}}(V)).$$
(155)

The proof is finished taking λ satisfying $\lambda^{-(\alpha+1)/\alpha} = t$.

Proof of Lemma 9.7. The claim of the lemma is equivalent with

$$\int_0^1 \bar{\mathbb{P}}(\bar{Z}(t) \notin J_{m'}^{\eta'}) dt \le \delta'. \tag{156}$$

We use $\sigma(m)$ to denote the first time \bar{Z} leaves [-m, m]. Let m' be large enough such that

$$\bar{\mathbb{P}}(\sigma(m') < 1) < \delta'/2. \tag{157}$$

and let $\sigma = \sigma(m')$. Then

$$\int_{0}^{1} \bar{\mathbb{P}}(\bar{Z}(t) \notin J_{m'}^{\eta'}) dt = \bar{\mathbb{E}} \Big[\int_{0}^{1} \mathbf{1} \{ \bar{Z}(t) \notin J_{m'}^{\eta'} \} dt \Big]
\leq \bar{\mathbb{E}} \Big[\int_{0}^{1} \mathbf{1} \{ \bar{Z}(t) \notin J_{m'}^{\eta'}, \sigma \geq 1 \} dt + \int_{0}^{1} \mathbf{1} \{ \sigma < 1 \} dt \Big]
\leq \bar{\mathbb{E}} \int_{0}^{\sigma} \mathbf{1} \{ \bar{Z}(t) \notin J_{m'}^{\eta'} \} dt + \delta'/2.$$
(158)

We should bound the expectation in the last expression by $\delta'/2$. We establish this bound by proving

$$\bar{\mathbb{P}}\Big[\bar{\mathbb{E}}\Big(\int_0^\sigma \mathbf{1}\{\bar{Z}(t) \notin J_{m'}^{\eta'}\} dt \Big| V\Big) \ge \delta'/4\Big] \le \delta'/4. \tag{159}$$

The conditional expectation inside the brackets can be written as

$$\bar{\mathbb{E}}\Big[\int_{0}^{\sigma} \mathbf{1}\{\bar{Z}(t) \notin J_{m'}^{\eta'}\} dt \Big| V\Big] = \sum_{\substack{x_{i} \in [-m', m'] \\ v_{i} < \eta'}} G_{m'}(0, x_{i}) v_{i}, \tag{160}$$

where as usually (x_i, v_i) is the collection of atoms of $\bar{\rho}$ and $G_m(x, y)$ is the Green's function of the standard Brownian motion killed on exit from [-m, m]. There exists a constant k depending only on m such that $G_m(0, x) \leq k$ for all $x \in [-m, m]$. We thus have

$$\bar{\mathbb{P}}\Big[\bar{\mathbb{E}}\Big(\int_0^\sigma \mathbf{1}\{\bar{Z}(t) \notin J_{m'}^{\eta'}\} dt \Big| V\Big) \ge \delta'/4\Big] \le \bar{\mathbb{P}}\Big[k \sum_{\substack{x_i \in [-m', m'] \\ v_i < \eta'}} v_i \ge \delta'/4\Big].$$
(161)

The sum in the last equation has the same distribution as the Lévy process V without jumps larger than η' at the time 2m'. One can thus easily choose η' small enough, such that the last probability is smaller than $\delta'/4$.

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